

At time t , the particle is at some point further along its path, but only the conditions at the time t_{past} determine the fields at point r at a time t . The magnetic field is always perpendicular to both \mathbf{n} and \mathbf{E}

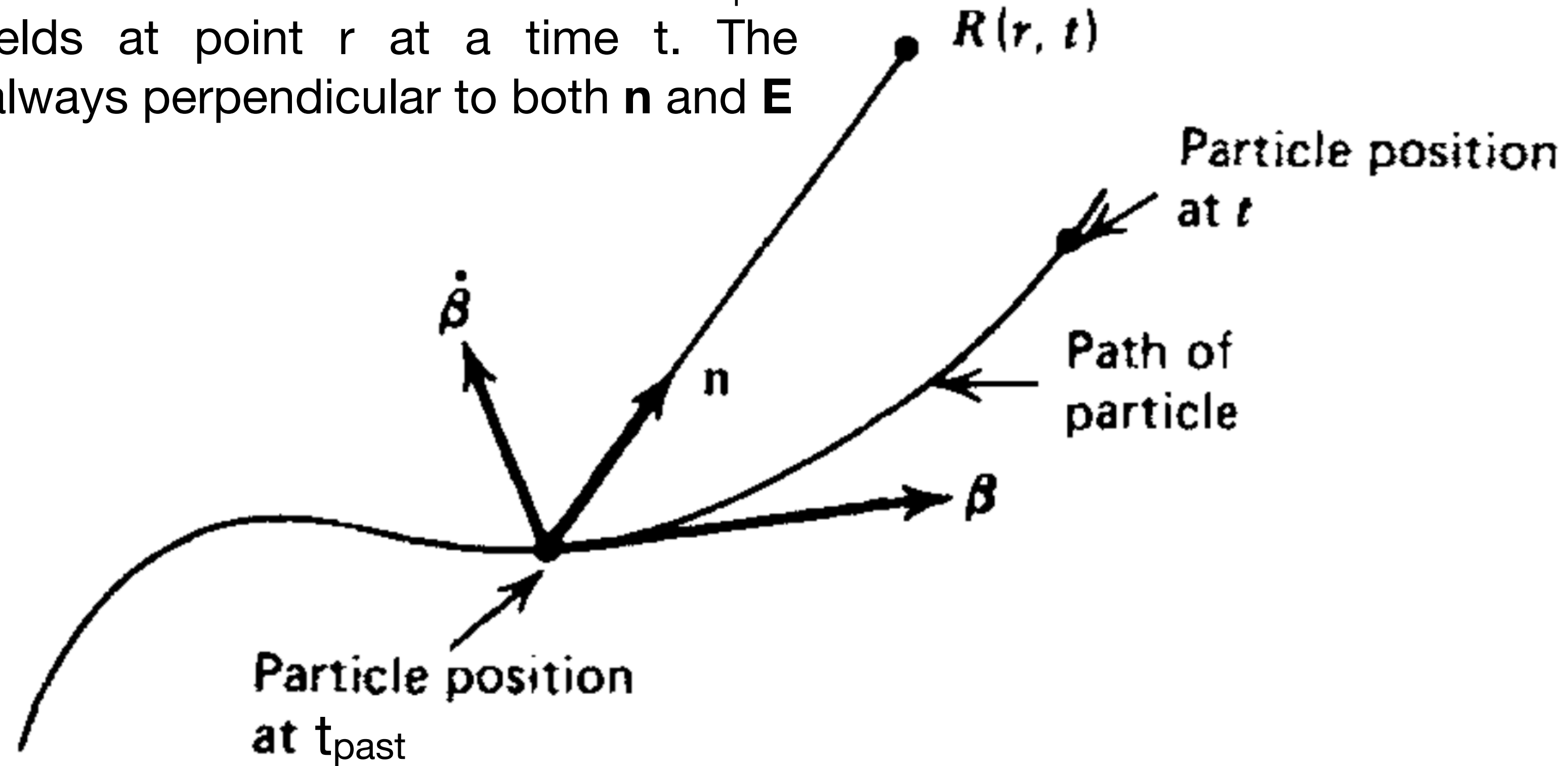


Figure 3.1 Geometry for calculation of the radiation field at R from the position of the radiating particle at the retarded time.

Drawn in the plane of \mathbf{n} and acceleration.

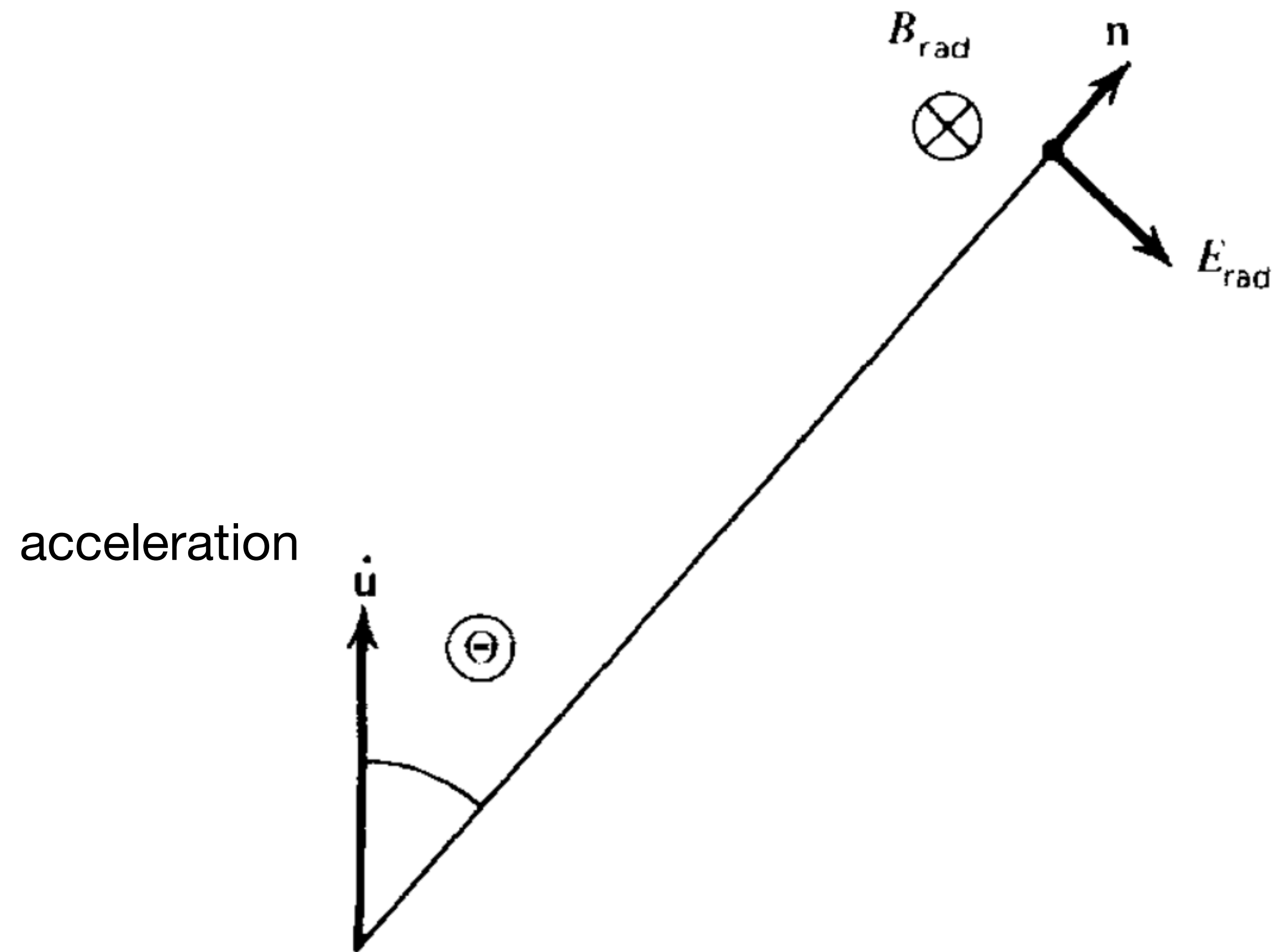


Figure 3.3 *Electric and magnetic radiation field configurations for a slowly moving particle. The direction of \mathbf{B}_{rad} is into the page.*

THOMSON SCATTERING

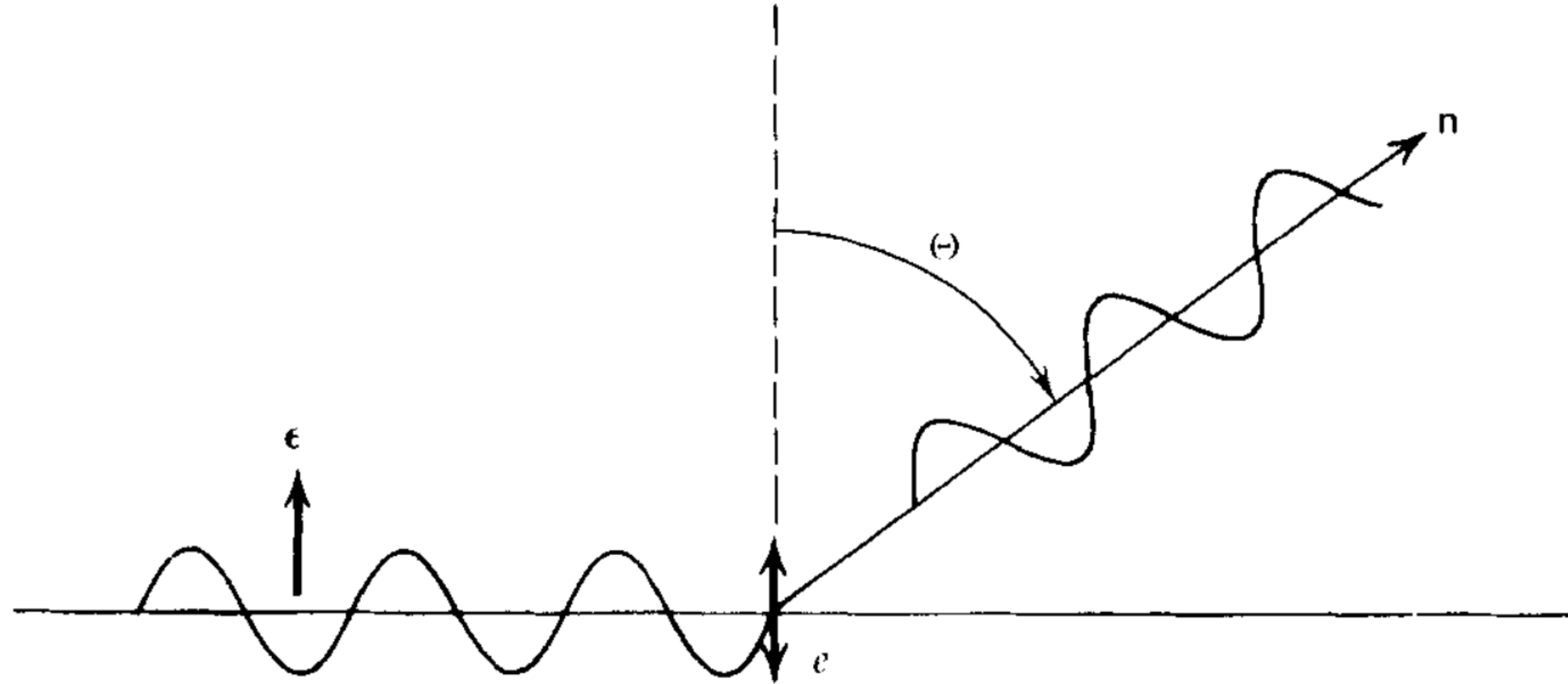


Figure 3.6 Scattering of polarized radiation by a charged particle.

The incident electromagnetic wave accelerates the electron, causing it to emit radiation at the same frequency of the incident wave, resulting in the wave being scattered.

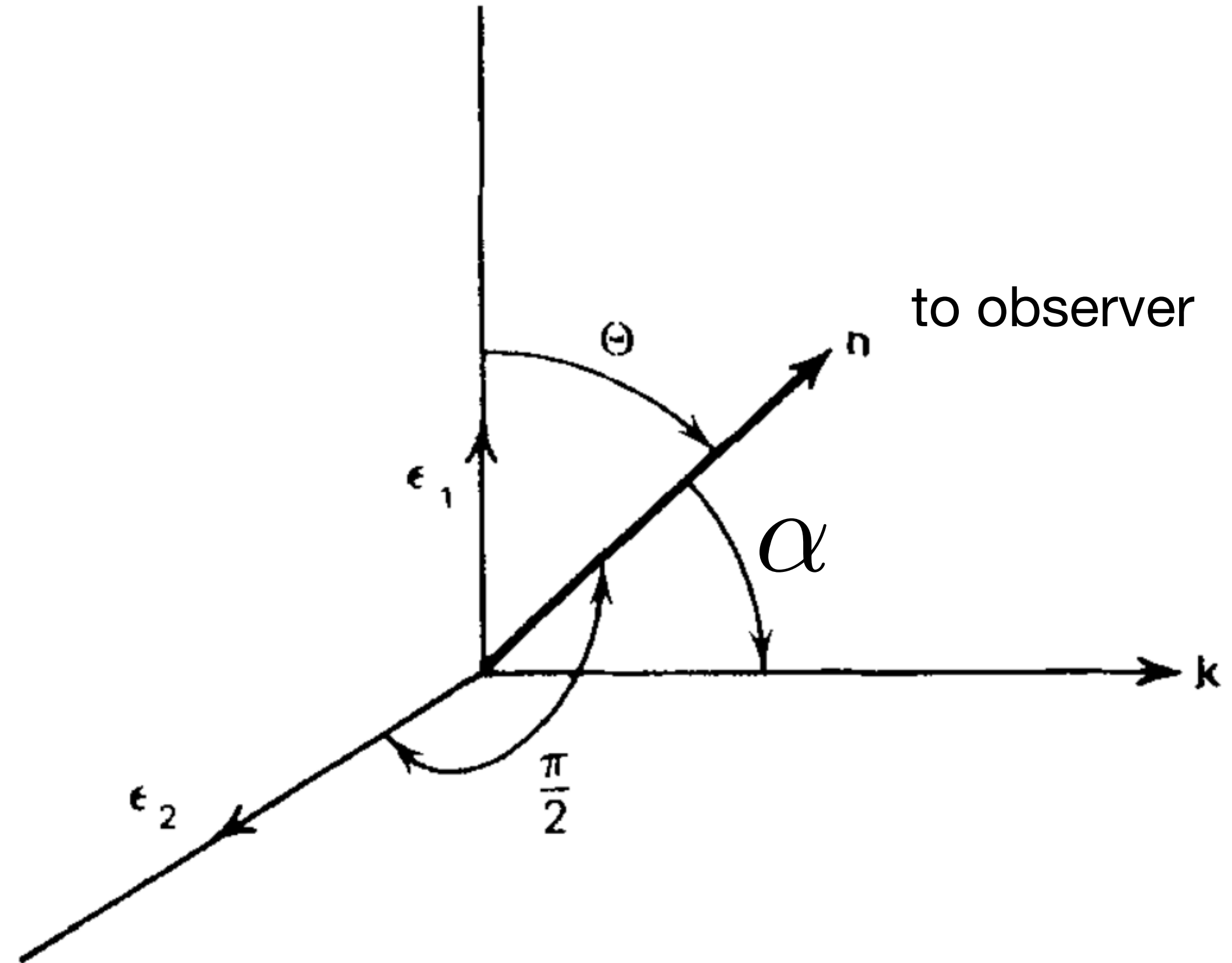


Figure 3.7 *Geometry for scattering unpolarized radiation.*

Approximate analytic formulas for $\langle g_{ff} \rangle$

U.P. = uncertainty principle dominated regime

$$u \equiv \frac{h\nu}{kT}$$

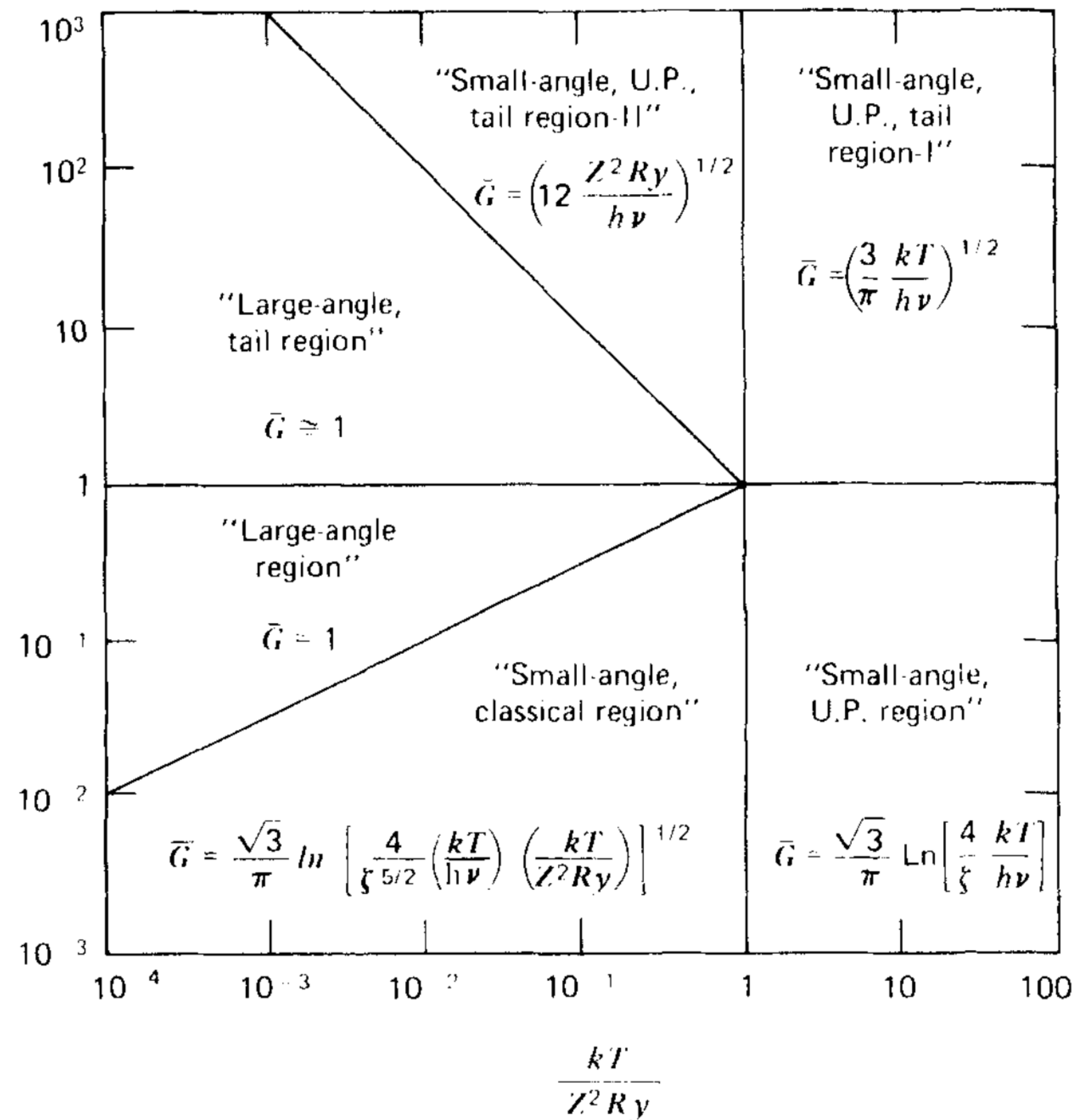


Figure 5.2 Approximate analytic formulae for the gaunt factor $\bar{g}_{ff}(\nu, T)$ for thermal bremsstrahlung. Here \bar{g}_{ff} is denoted by \bar{G} and the energy unit $Ry = 13.6$ eV. (Taken from Novikov, I. D. and Thorne, K. S. 1973 in *Black Holes, Les Houches*, Eds. C. DeWitt and B. DeWitt, Gordon and Breach, New York.)

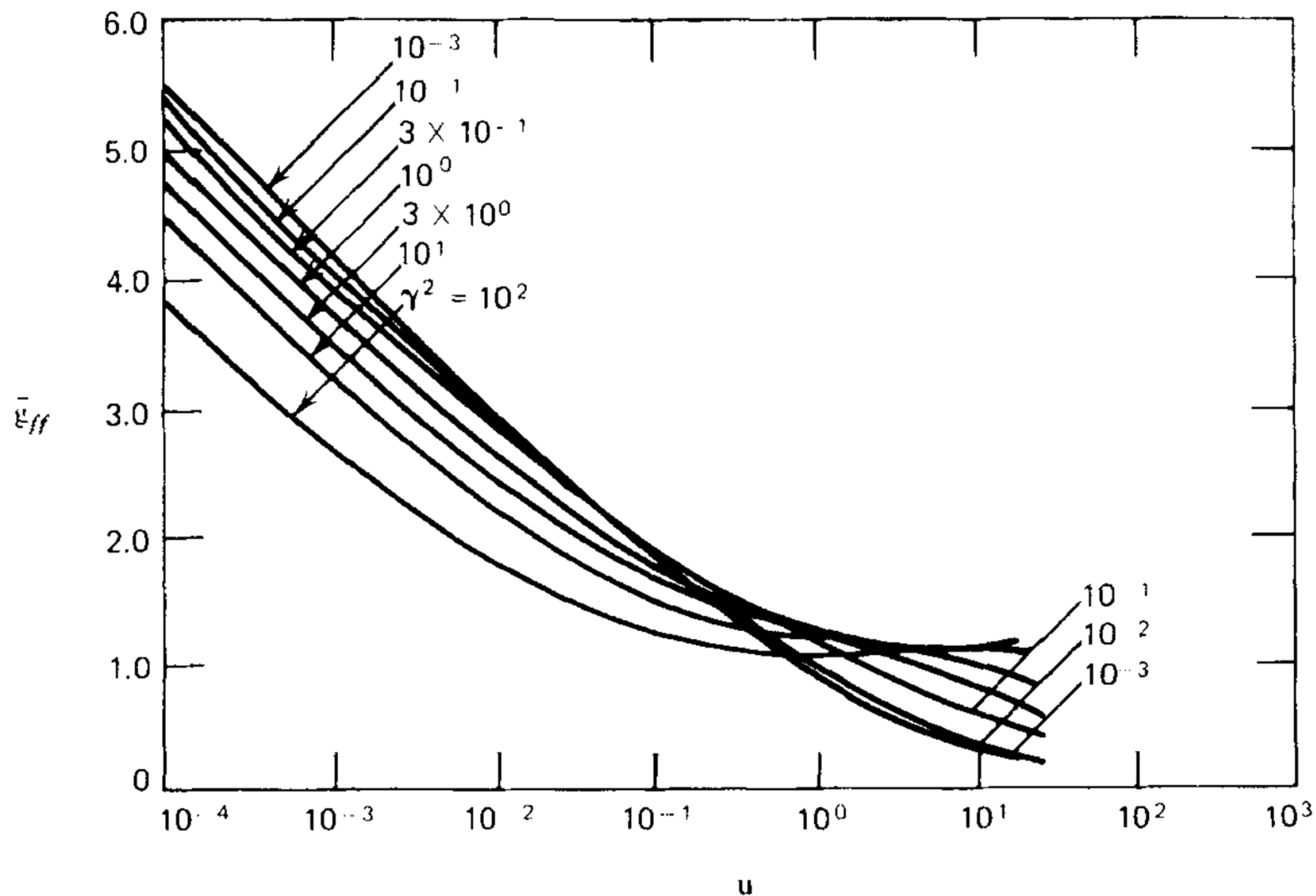


Figure 5.3 Numerical values of the gaunt factor $\bar{g}_{ff}(\nu, T)$. Here the frequency variable is $u = 4.8 \times 10^{11} \nu / T$ and the temperature variable is $\gamma^2 = 1.58 \times 10^5 Z^2 / T$. (Taken from Karzas, W. and Latter, R. 1961, *Astrophys. J. Suppl.*, 6, 167.)

- Spectrum of Thermal Bremsstrahlung

