

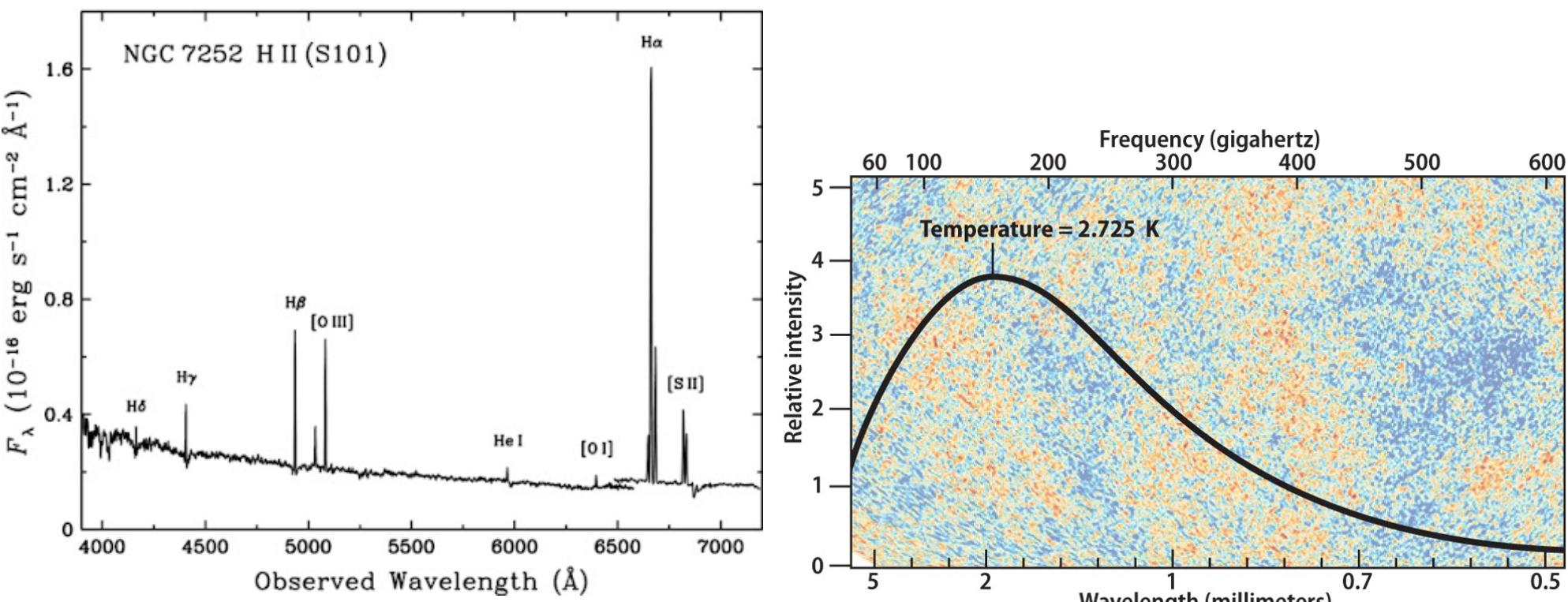
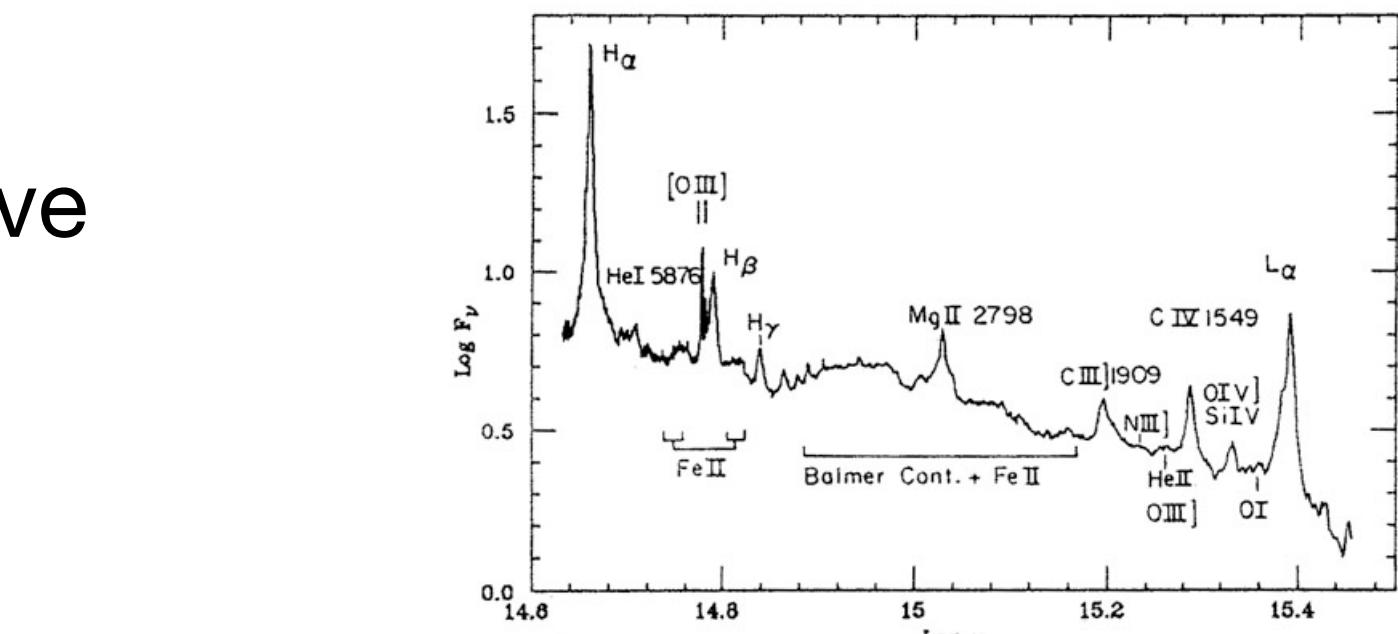
AST-191 Special Topics: “Radiative Processes in Astrophysics”

Prof. Danilo Marchesini
danilo.marchesini@tufts.edu

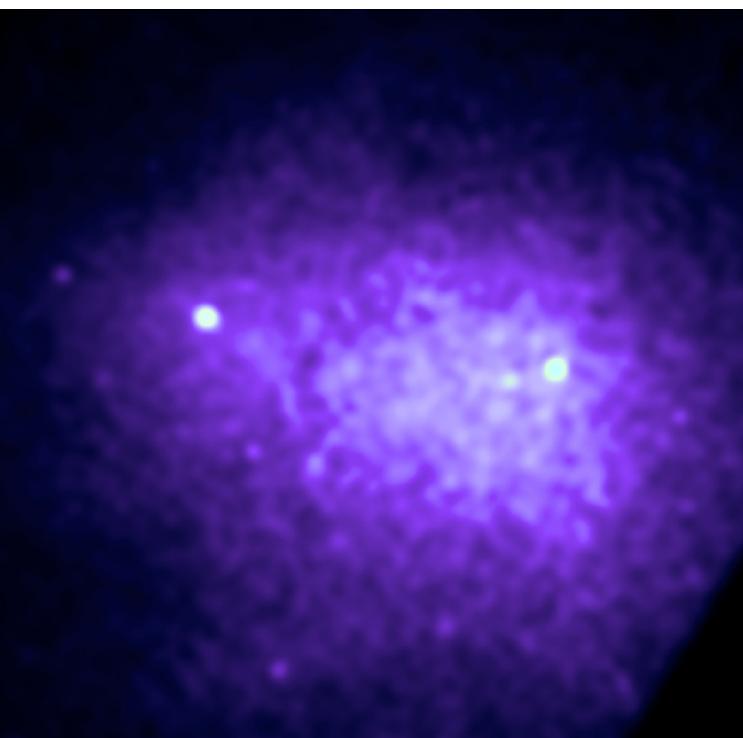
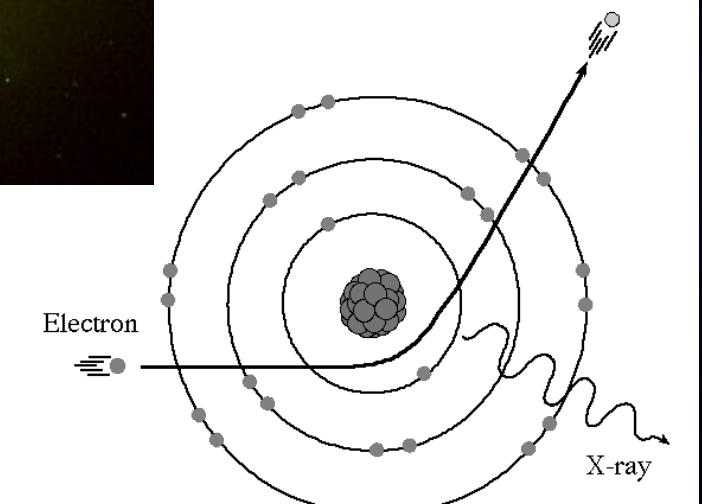
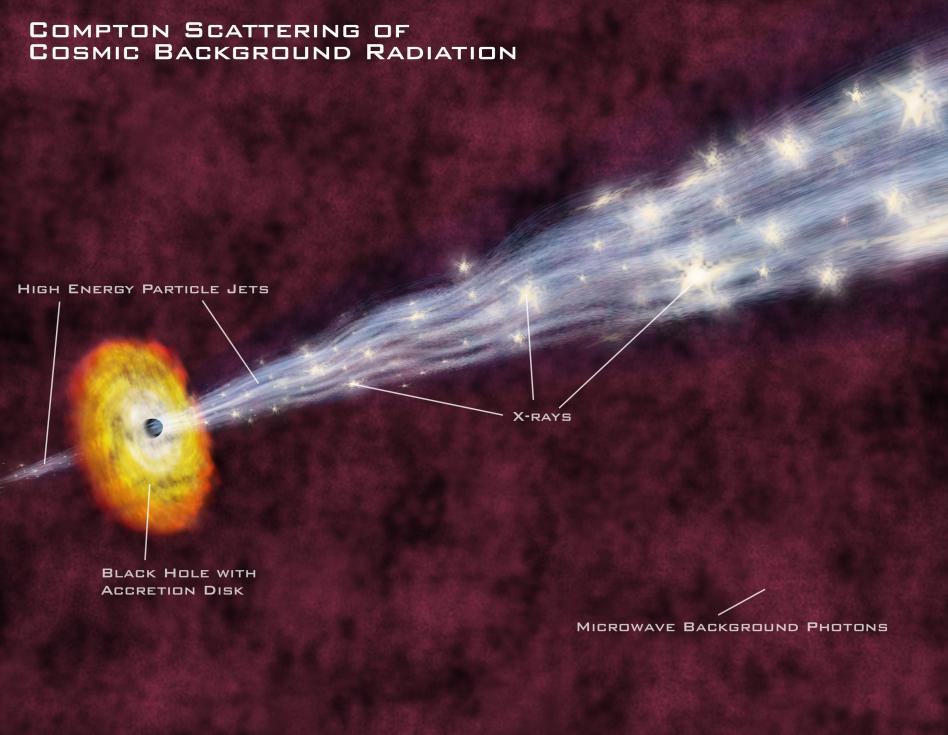
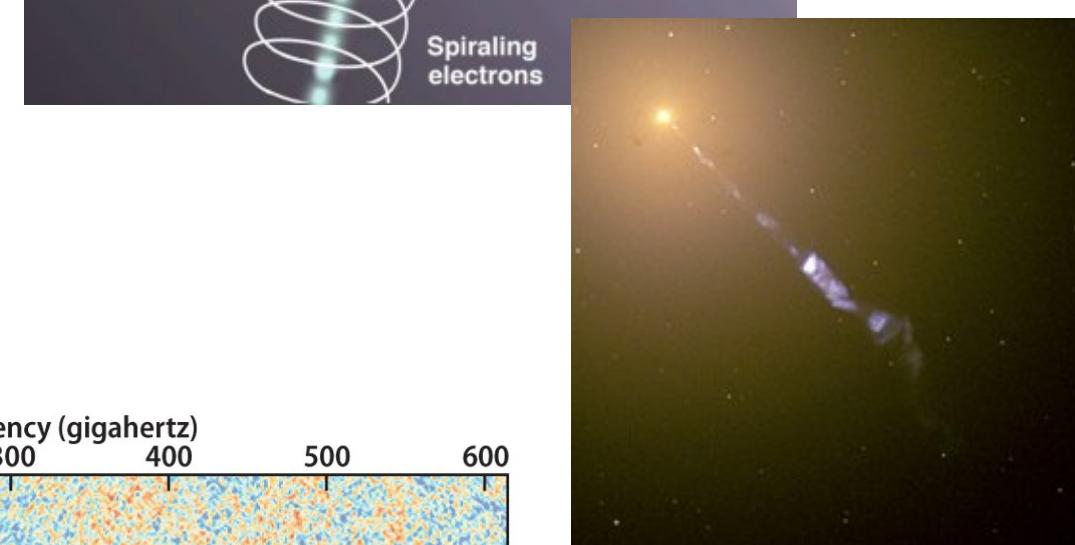
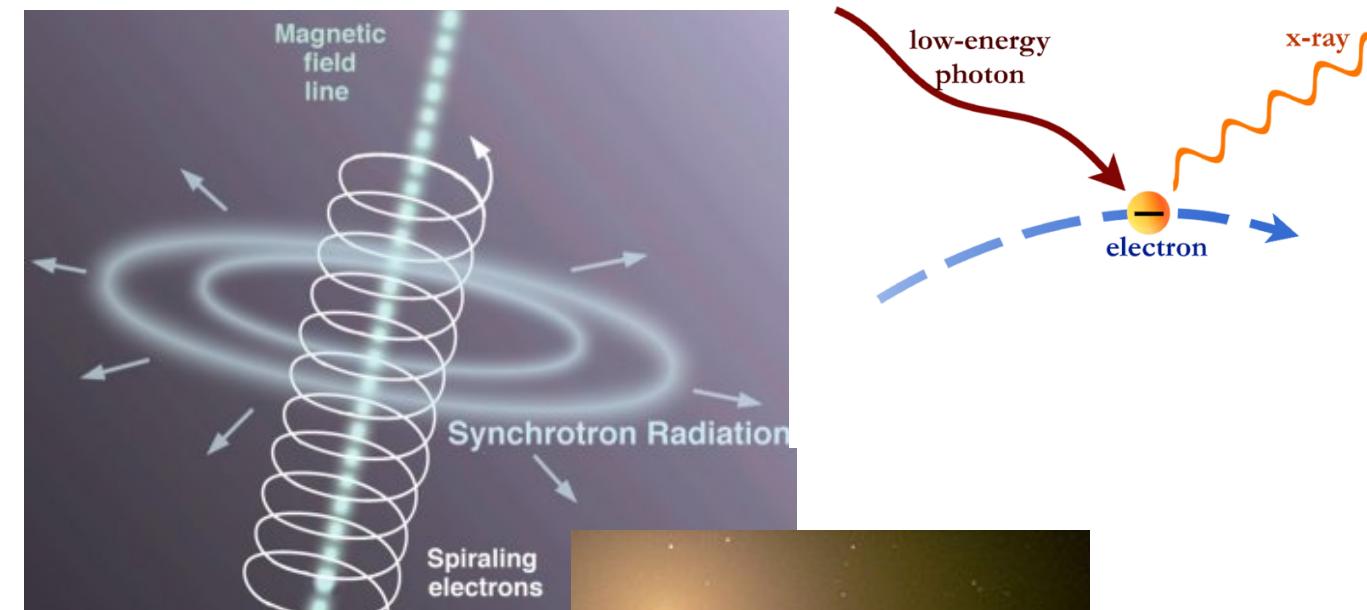
Fall 2024 semester:
 Mon-Wed
 9:00-10:15am
 CLIC 204

COVERED TOPICS:

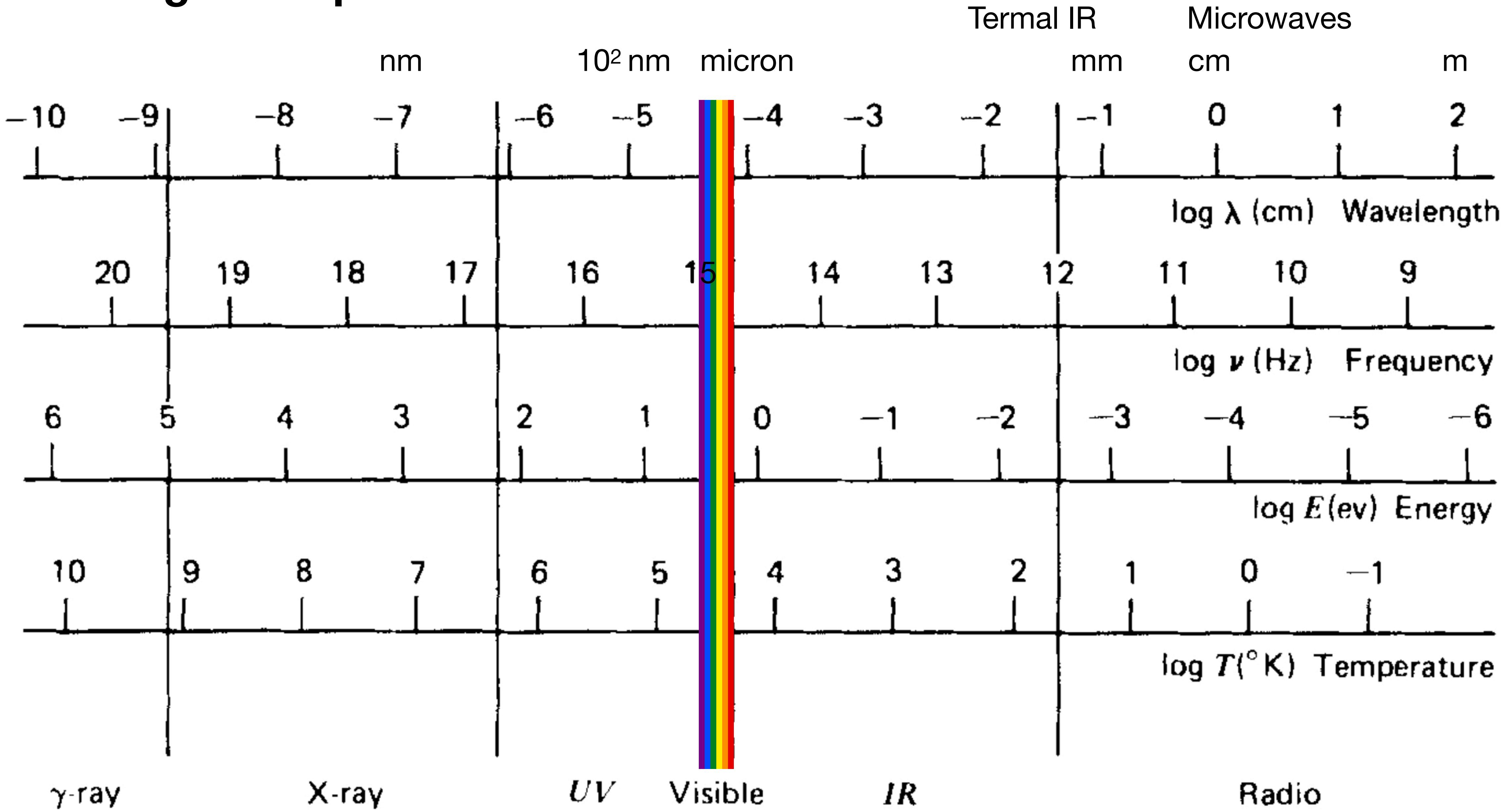
- Fundamentals of Radiative Transfer
- Radiation from Moving Charges
- Bremsstrahlung
- Synchrotron Radiation
- Compton scattering
- Plasma Effects
- Atomic Structure
- Radiative Transitions
- Molecular Structure

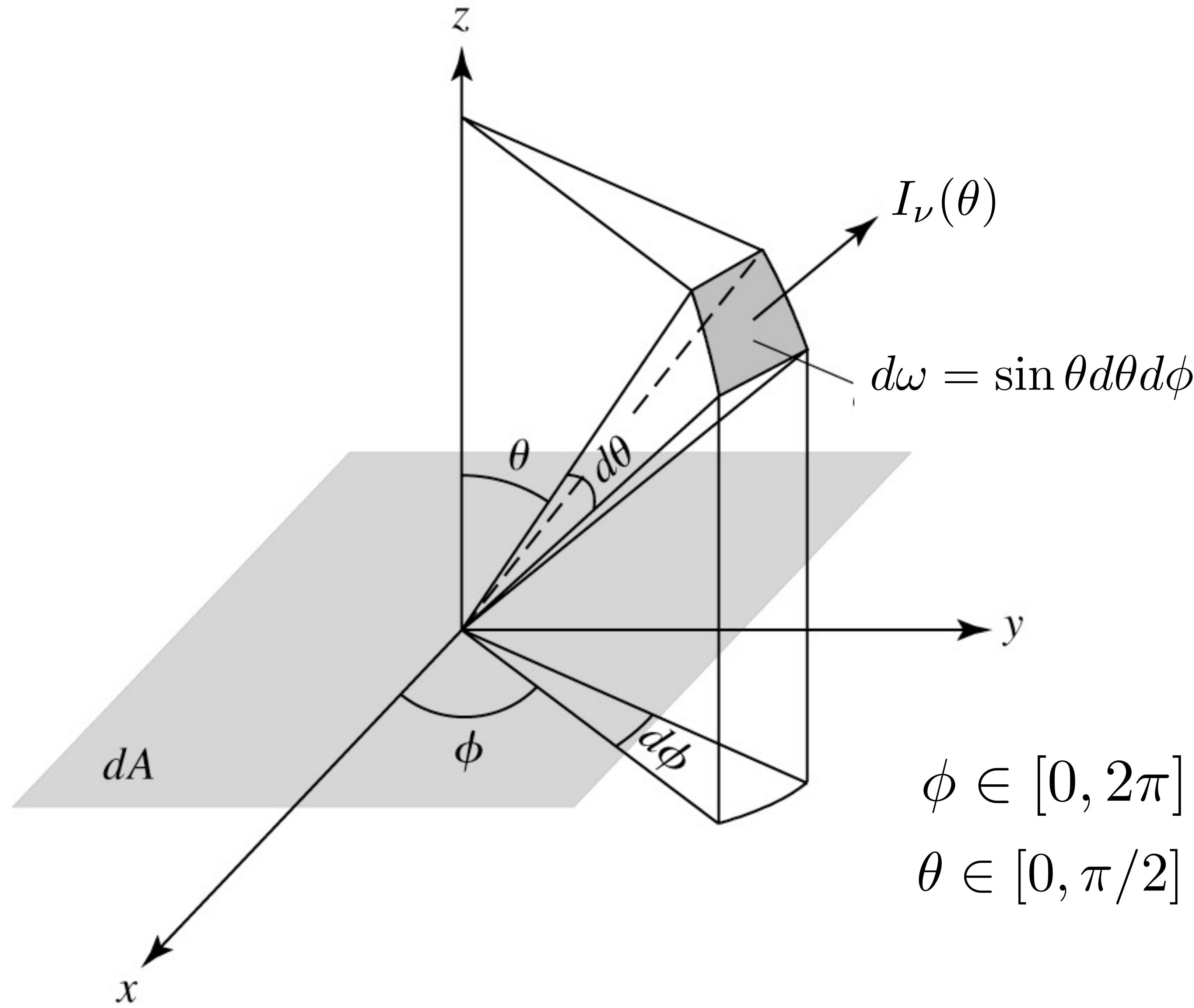


$\lambda\nu = c$	$E = h\nu = hc/\lambda$	$P = E/c$	$W = 2\pi\nu$	$P = \hbar k$	$m - m_0 = -2.5 \log(F/F_0)$	$L = \frac{dE}{dt}$
ELUX $F_\nu = dE/dAdt+d\nu$	TOTAL FLUX $F = \int F_\nu d\nu$				INVERSE SQUARE LAW $F = L/4\pi r^2$	
BRIGHTNESS OR SPECIFIC INTENSITY $I_\nu = \frac{dE}{dAd\Omega dt d\nu} \left[\frac{\text{erg}}{\text{cm}^2 \text{ster Hz}} \right]$	$F_\nu = \int I_\nu \cos\theta d\Omega \left[\frac{\text{erg}}{\text{cm}^2 \text{Hz}} \right]$	$F \left[\frac{\text{erg}}{\text{cm}^2} \right]$				
ENERGY DENSITY $U_\nu = \frac{1}{c} \int I_\nu d\Omega = \frac{4\pi}{c} J_\nu$	$J_\nu = \frac{1}{4\pi} \int I_\nu d\Omega$	$U = \int U_\nu d\nu = \frac{4\pi}{c} \int J_\nu d\nu$				
RADIATION PRESSURE (ISOTROPIC) $P_\nu = \frac{2}{c} \int I_\nu \cos^2\theta d\Omega$	$P = \frac{2}{c} \int J_\nu d\nu \int \cos^2\theta d\Omega$	$P = \frac{1}{3} u$	KIRCHHOFF'S LAW	$j_\nu = d\nu B_\nu$		
CONSTANCY OF I_ν $dE = I_\nu dA dt d\Omega d\nu$, $dA d\Omega = dAd\Omega d\nu$, $dE = dE_2 \Rightarrow I_\nu = I_\nu 2$	$I_\nu = I_\nu 2$	$I_\nu = \text{CONST.}$				
IN FREE SPACE $dE_2 = I_\nu dA_2 dt d\Omega_2 d\nu_2$, $d\nu_2 = d\nu$	$\frac{dI_\nu}{ds} = 0$	$dI_\nu/ds = 0$				
UNIFORMLY BRIGHT SPHERE $I_\nu = B_\nu$	$F = \int I_\nu \cos\theta d\Omega = B_\nu \int \sin\theta d\Omega \int \cos\theta d\Omega$	$F = \pi B \left(\frac{R}{r} \right)^2$	$F = \pi B$	$I(\theta) = \begin{cases} B & \theta = 0 \\ 0 & \theta > 0 \end{cases}$		
$\theta_c = \sin^{-1}(R/r)$	$\Rightarrow F = \pi B (1 - \cos^2\theta_c) = \pi B \sin^2\theta_c$	$\theta = R/r$	$r = R$			
RADIATIVE TRANSFER EMISSION COEFFICIENT $j_\nu = \frac{dE}{dV dt d\nu d\Omega}$	$dI_\nu = j_\nu ds$	$\frac{dI_\nu}{ds} = j_\nu \neq 0$	ISOTROPIC Emitter $j_\nu = \frac{1}{4\pi} P_\nu = \frac{E_\nu P}{4\pi r^2}$	$P = \frac{\text{Power}}{\text{Area}}$		
ABSORPTION $dI_\nu = -\alpha_\nu I_\nu ds = -n_\nu \sigma_\nu I_\nu ds$	$\alpha_\nu = n_\nu \sigma_\nu = QK_\nu$	$\alpha_\nu [\text{cm}^{-1}]$	$\sigma_\nu [\text{cm}^{-2}]$ CROSS SECTION	$K_\nu [\text{cm}^2 \text{g}^{-1}]$ OPACITY		
RADIATIVE TRANSFER EQ $\frac{dI_\nu}{ds} = -\alpha_\nu I_\nu + j_\nu$	EMISSION ONLY $\frac{dI_\nu}{ds} = j_\nu \quad I_\nu(s) = I_\nu(s_0) + \int_{s_0}^s j_\nu(s') ds'$	$\frac{dI_\nu}{ds} = j_\nu \quad I_\nu(s) = I_\nu(s_0) e^{-\int_{s_0}^s \alpha_\nu(s') ds'}$	ABSORPTION ONLY $\frac{dI_\nu}{ds} = -I_\nu + S_\nu$	$S_\nu = \frac{j_\nu}{\alpha_\nu}$	$I_\nu(s) = I_\nu(s_0) e^{-\int_{s_0}^s \alpha_\nu(s') ds'}$	$I_\nu(s) = I_\nu(s_0) e^{-\int_{s_0}^s \alpha_\nu(s') ds'}$
OPTICAL DEPTH $d\tau_\nu = d\alpha_\nu ds$	$\tau_\nu(s) = \int_{s_0}^s \alpha_\nu(s') ds'$	$\tau_\nu > 1$ OPTICALLY THICK $\tau_\nu < 1$ OPTICALLY THIN	$\frac{dI_\nu}{ds} = -I_\nu + S_\nu$	$S_\nu = \frac{j_\nu}{\alpha_\nu}$	$I_\nu(s) = I_\nu(s_0) e^{-\int_{s_0}^s \alpha_\nu(s') ds'}$	$I_\nu(s) = I_\nu(s_0) e^{-\int_{s_0}^s \alpha_\nu(s') ds'}$

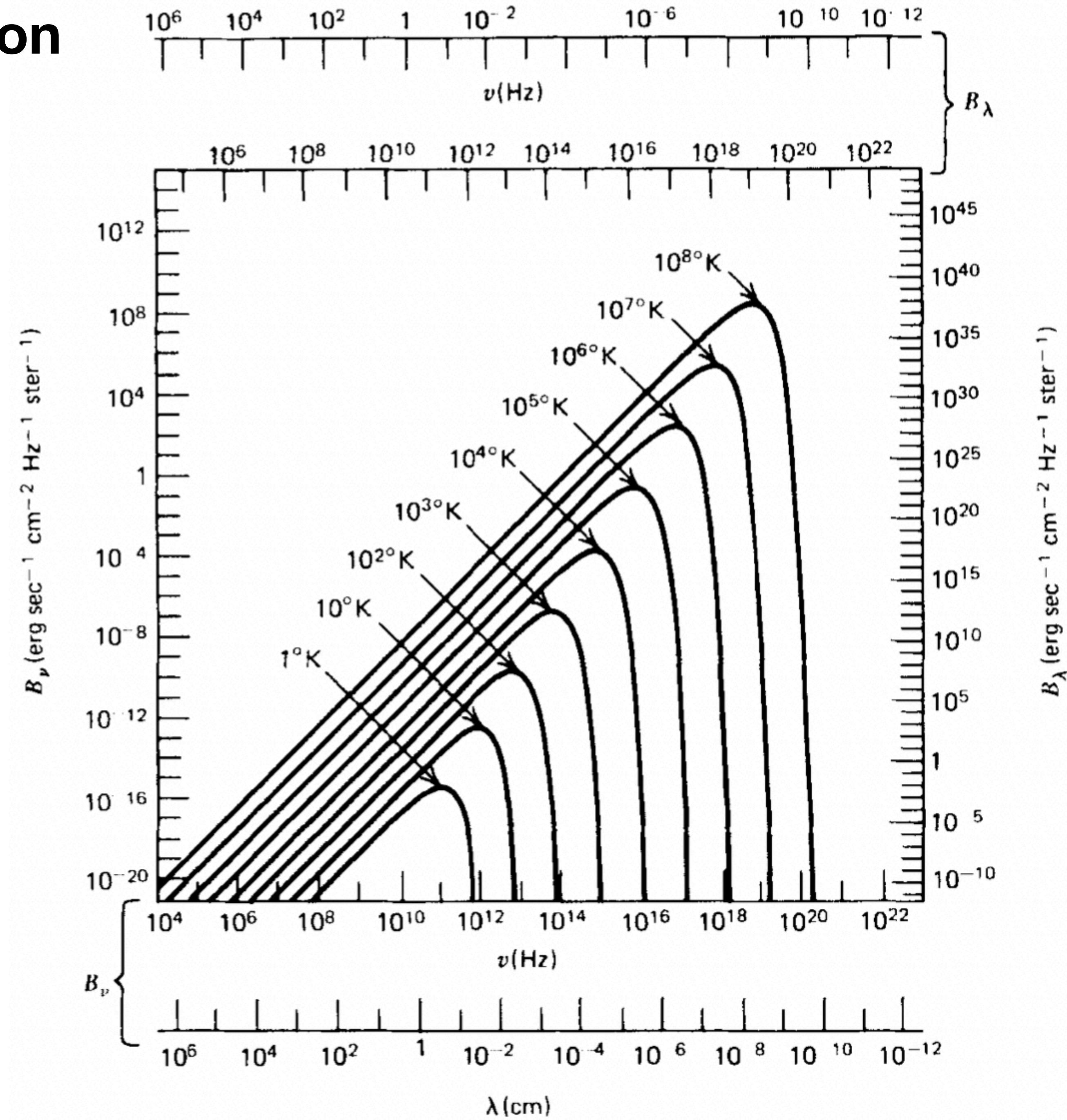


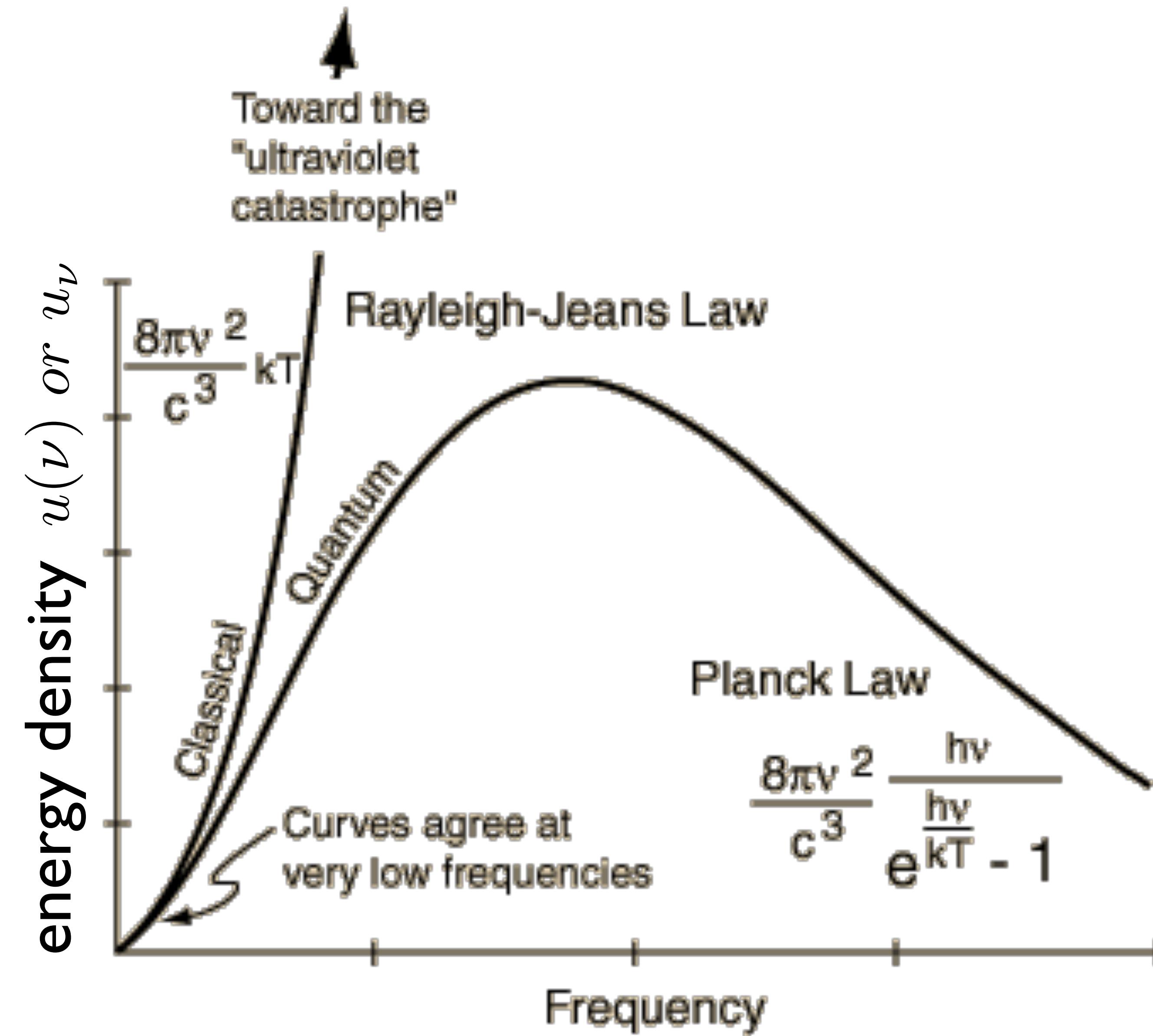
Electro-magnetic spectrum





Blackbody radiation (Plank's law)





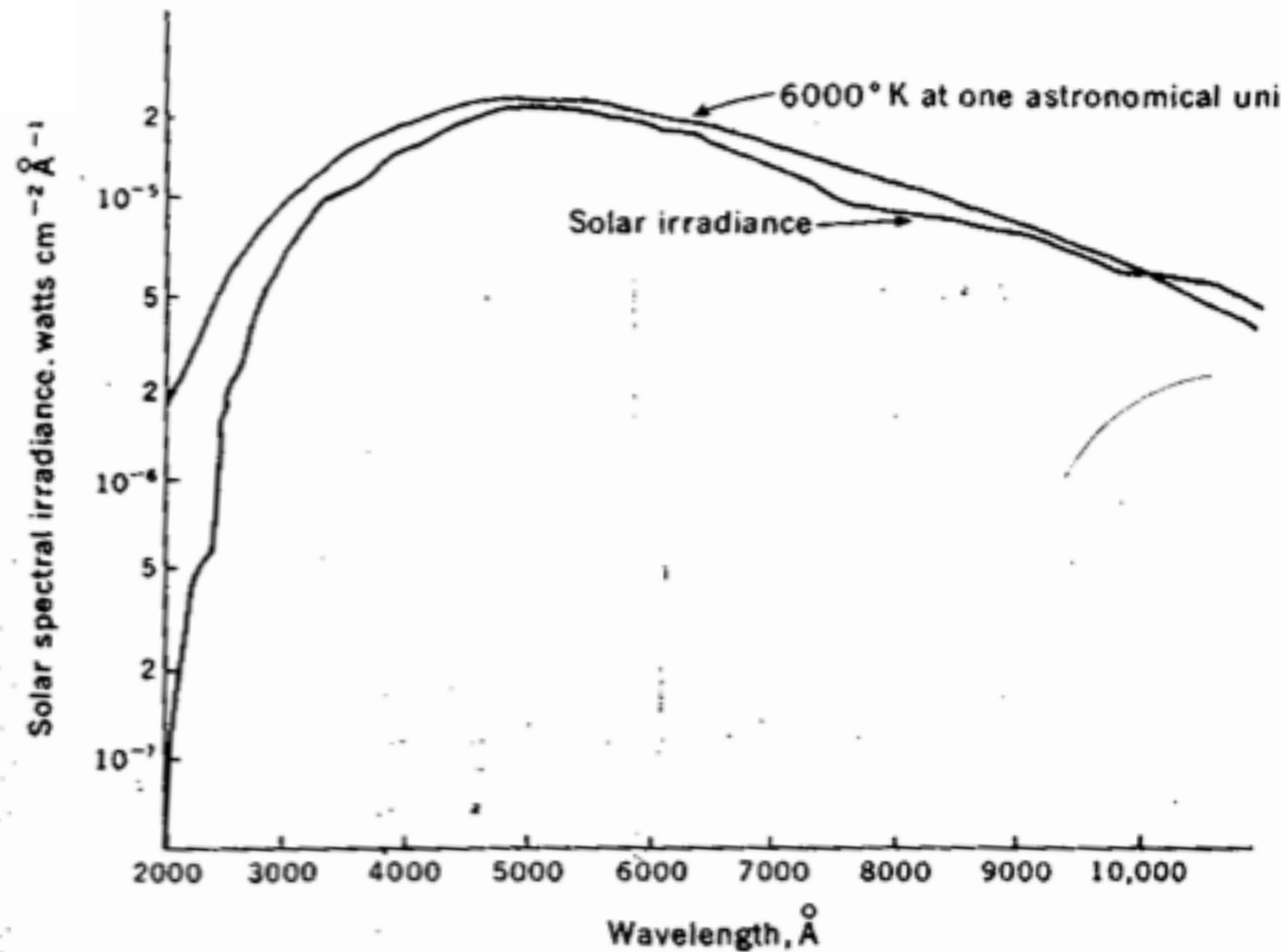
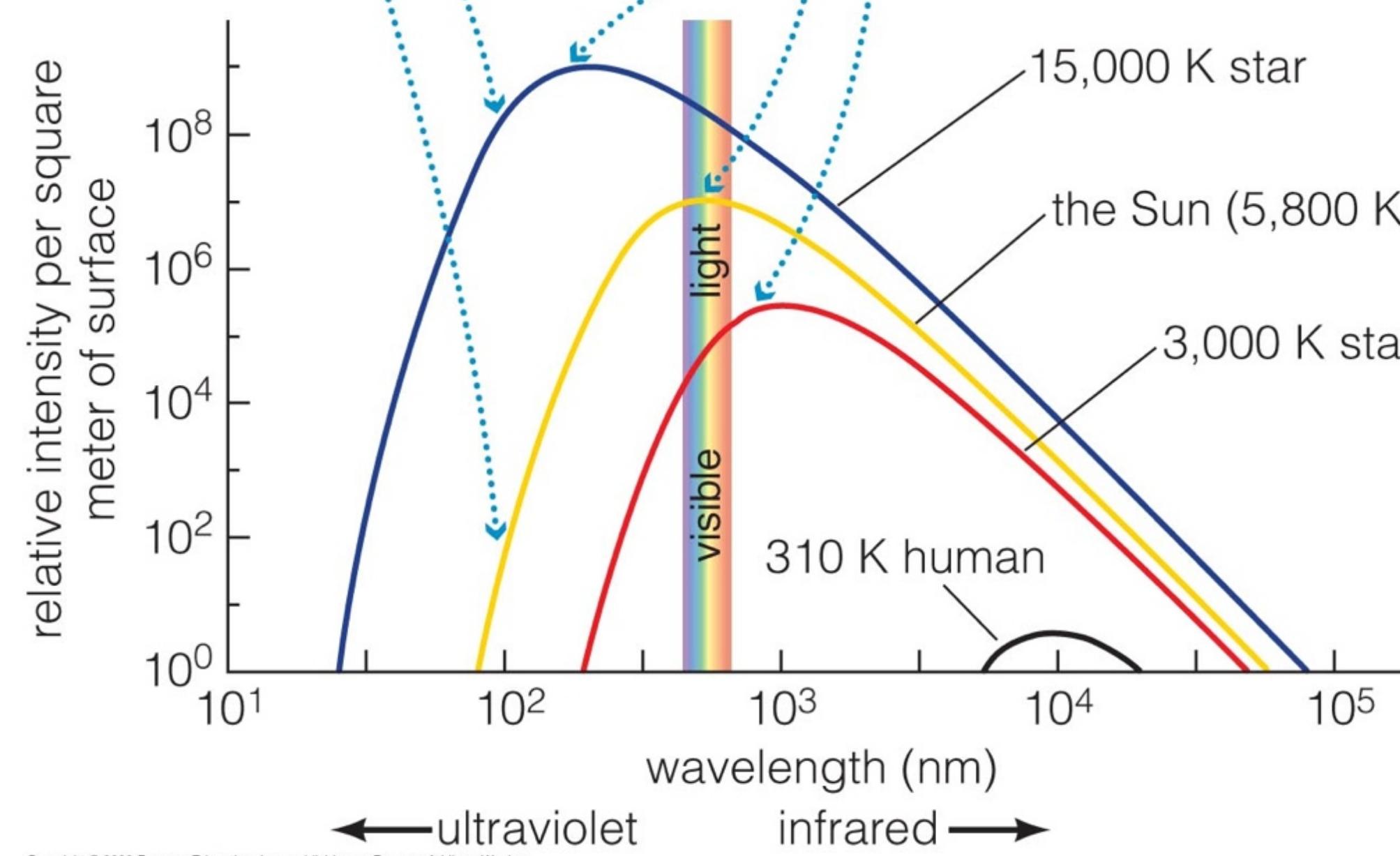


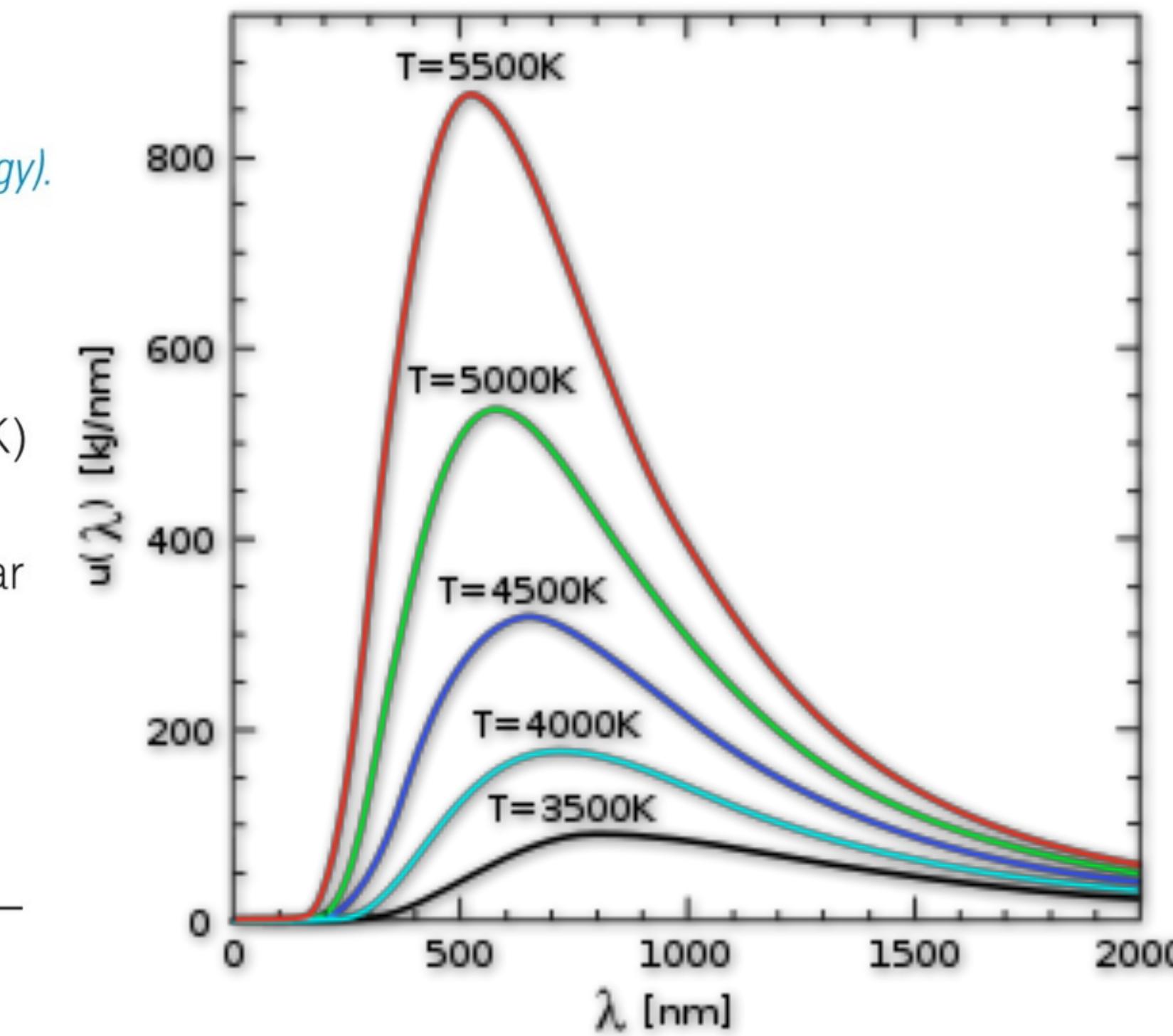
Fig. 1-5 Comparison of the visible solar energy-distribution curve with that from a blackbody at 6000°K. The overall resemblance is good, although the sun is quite deficient in the ultraviolet. [D. P. Le Galley and A. Rosen (eds.), "Space Physics," p. 111, John Wiley & Sons, Inc., New York, 1964.]

The curve for a hotter object is everywhere above the curve for a cooler object, showing that hotter objects emit more radiation per unit surface area at every wavelength.

The peak wavelength is further to the left for hotter objects, showing that hotter objects emit more of their light at shorter wavelength (high energy).

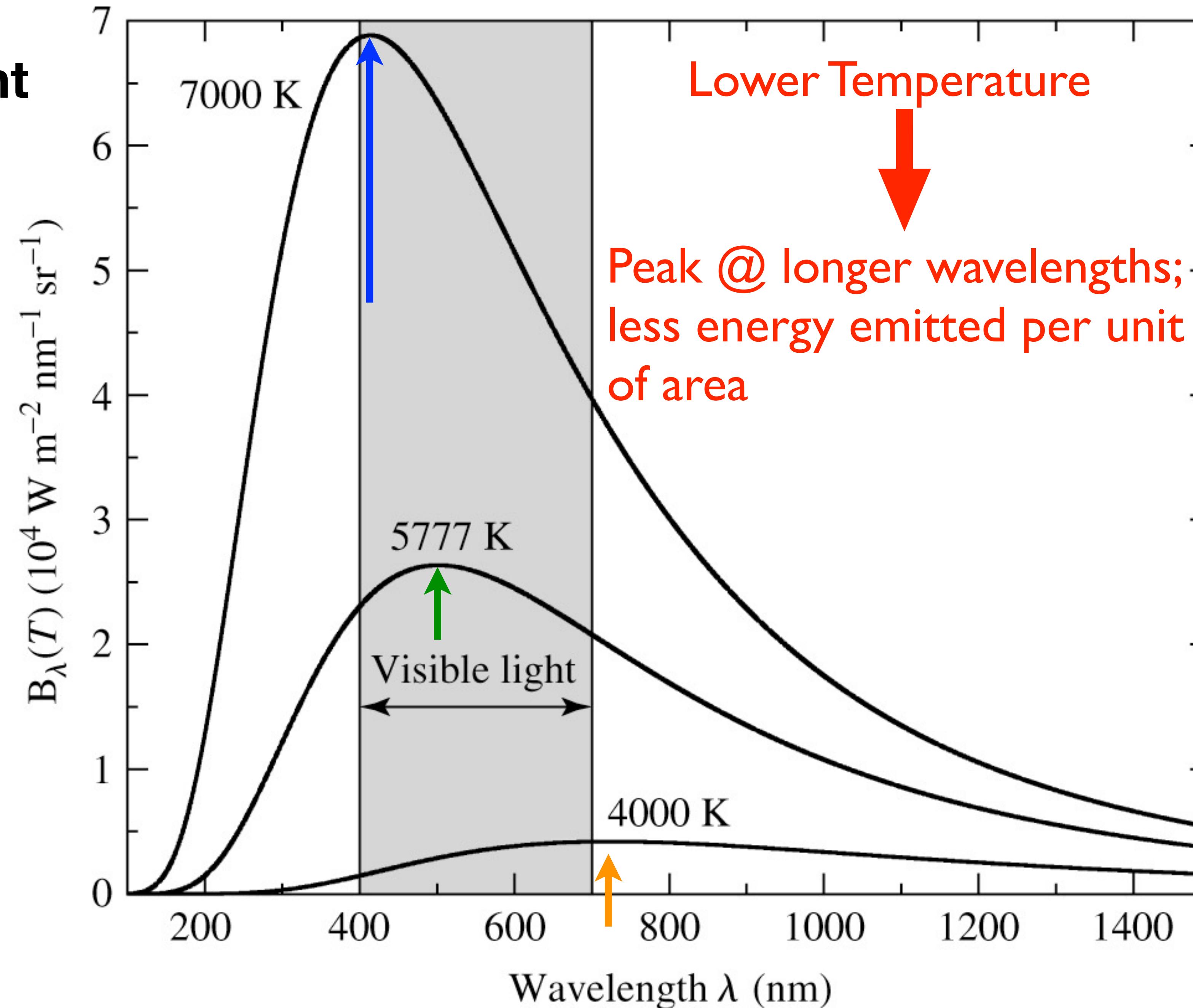


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Wien's displacement law

Wien's displacement law



Color temperature T_c

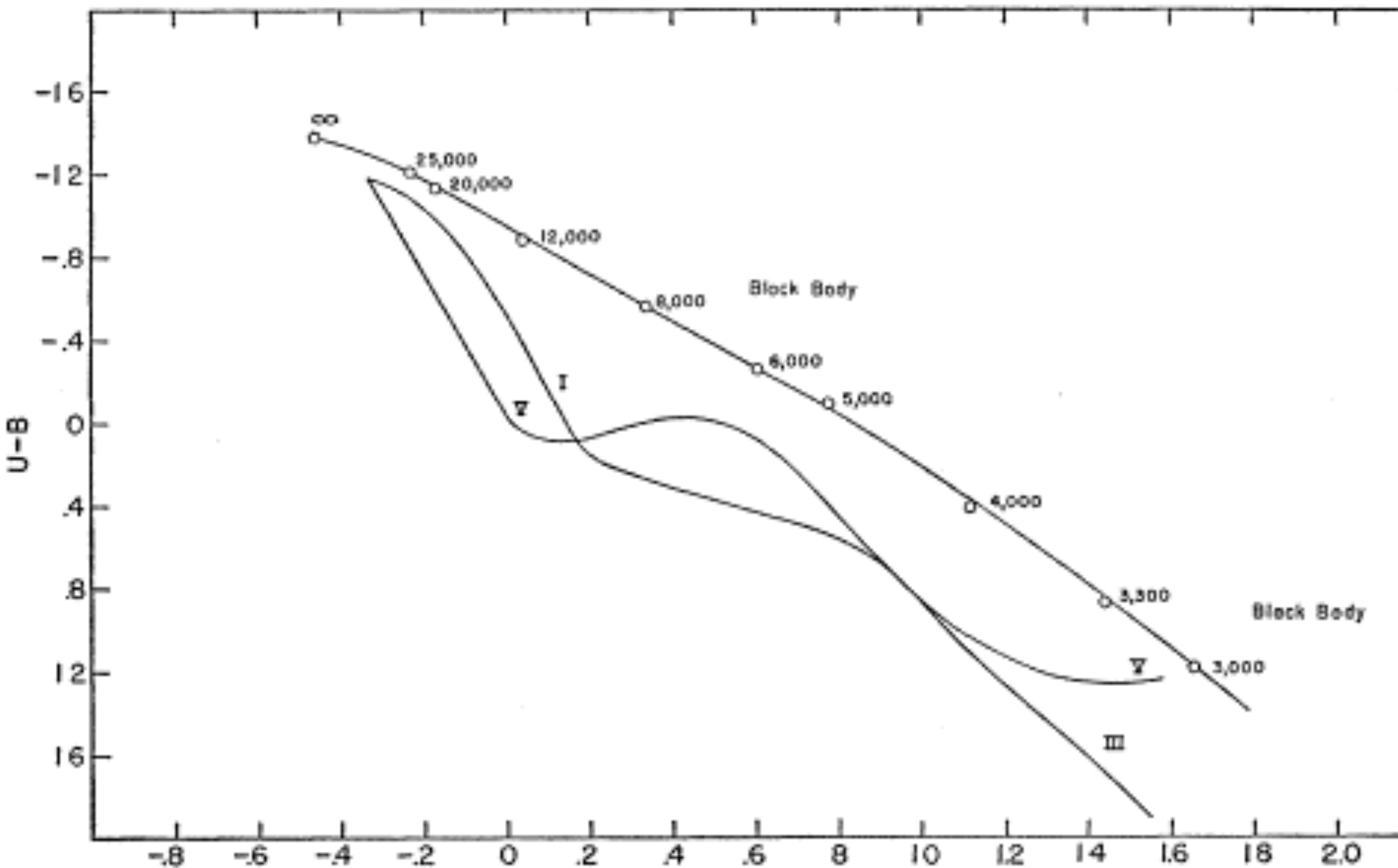
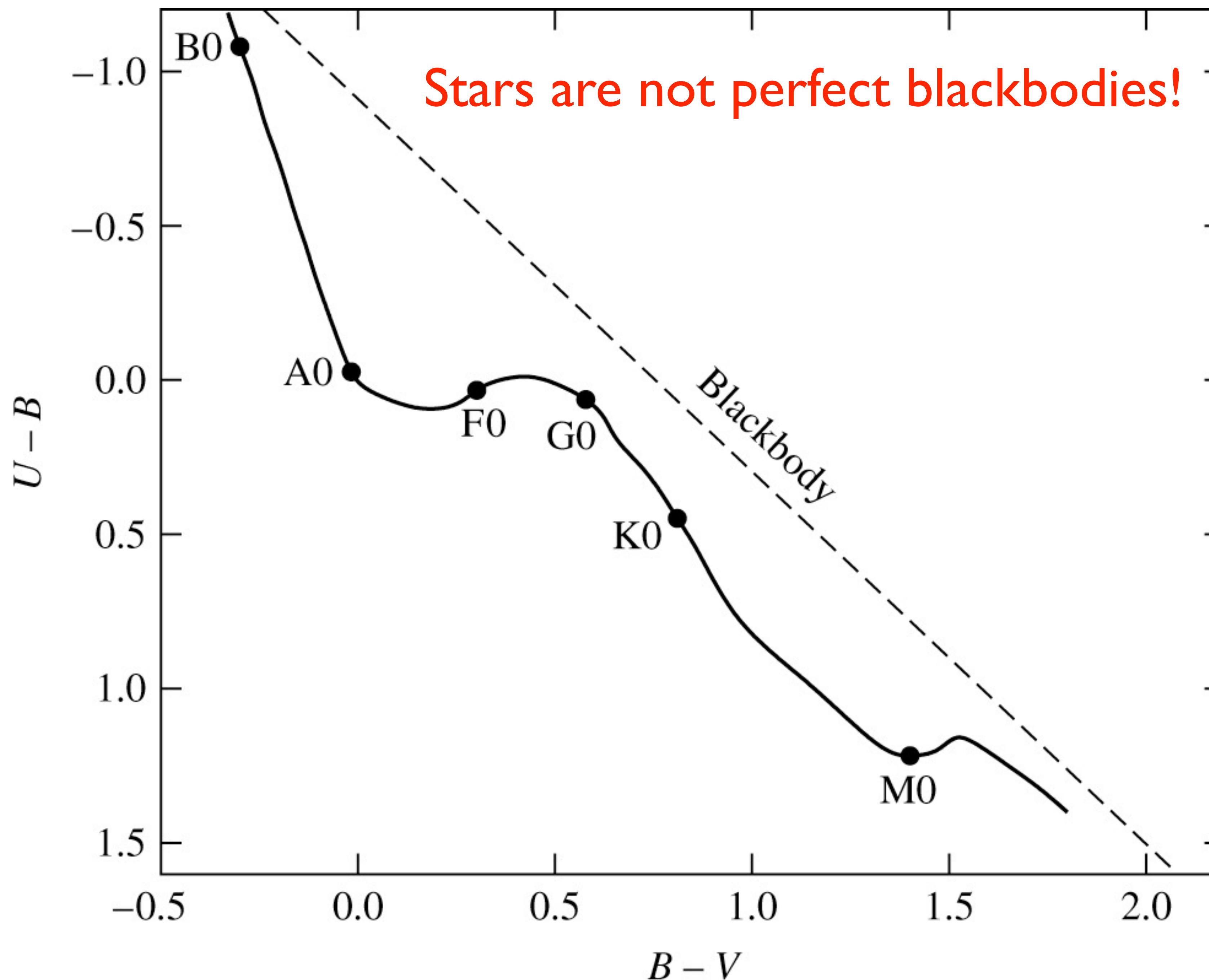
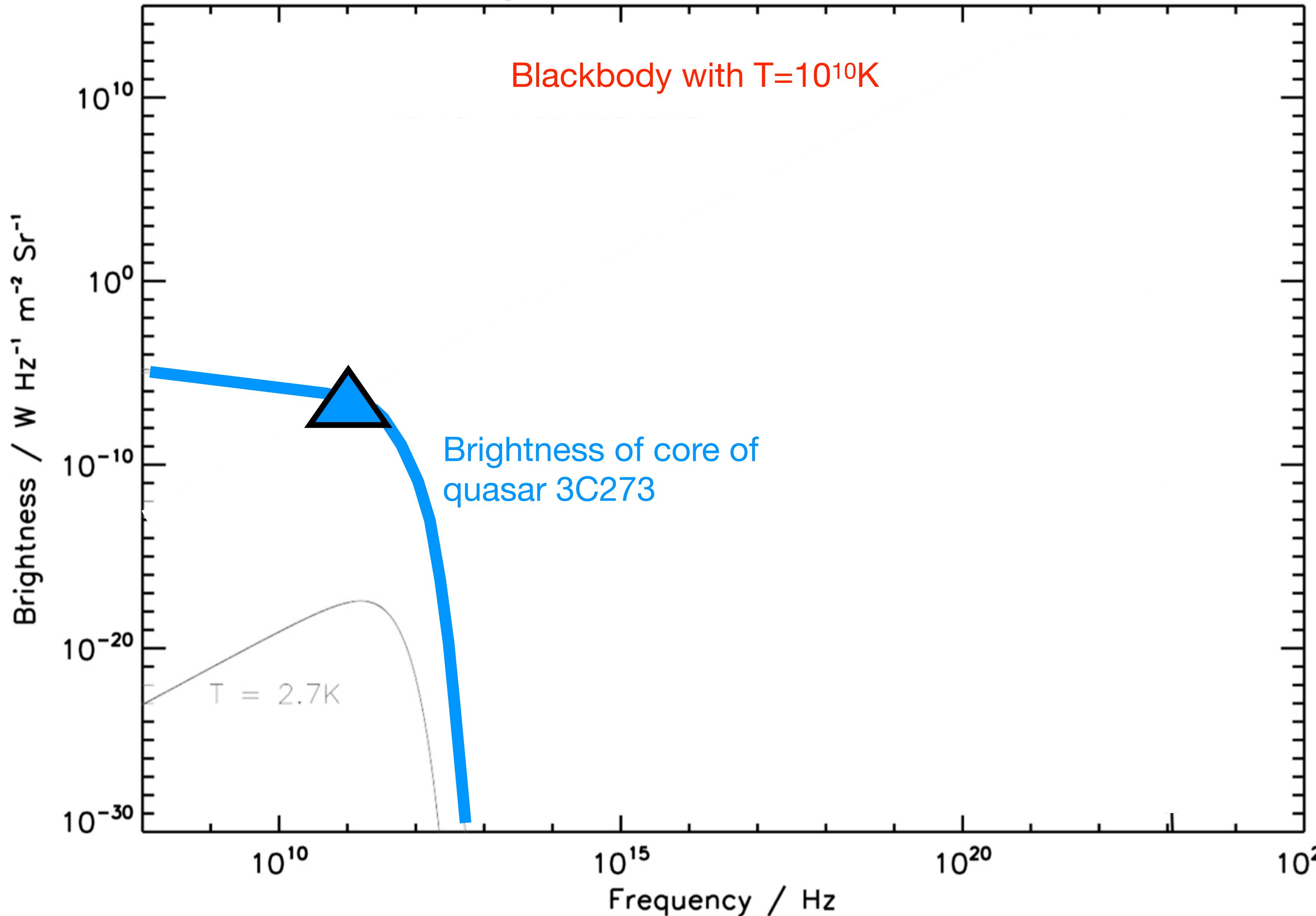


FIG. 1.—The $U - B$, $B - V$ relation for black bodies. The plotted points are taken from the last two columns of Table 2. The standard Johnson relations for stars of luminosity classes I–V are also shown.



Equivalent (brightness) temperature of a radiosource



CAREFUL! If the emitted spectrum is not blackbody radiation, then temperature measurements of T_b , T_c , or T_{eff} are wrong.