Cosmic Strings post Planck Adam Moss

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UNITED KINGDOM · CHINA · MALAYSIA



Observables



- CMB power spectrum. Includes recombination and postrecombination physics. Strongest limit on allowed fraction of strings, improved by Planck
- Non Gaussianity in CMB maps. Search for signatures of postrecombination Doppler shift induced by moving strings. Strongest limit also from Planck (see talk by Paul Shellard)
- CMB B modes. Defects produce comparable scalar, vector and tensor fluctuations. (see talk by Robert Brandenberger)
- 21 cm (see talk by Robert Brandenberger)
- Gravitational lensing (not discussed here)
- Pulsar timing/gravitational waves. Stochastic background from loops, waves from cusps (see talk by Richard Battye)



String Evolution



- Perturbations from strings are active, so their evolution is key to understanding CMB anisotropies
- Strings evolve toward self-similar scaling regime
- Average properties of network are (nearly) constant with time
- Dynamics can be studied using numerical simulations
- Two approaches Nambu and Abelian-Higgs models
- Both have advantages and disadvantages
- Main issue is dynamical range assumptions have to be made in either case
- Will present Planck constraints for each case



Nambu Model

planck

- Thin string approximation
- Ignore radiation back reaction
- Impose reconnection by hand
- Network characterised by correlation length L
- Energy density is $\rho = \frac{\mu}{L^2}$,
- Observationally, string tension µ is main quantity of interest
- ► Find scaling solution L ~ t
- Make measurements of correlation length, velocity, small scale structure (wiggliness)



Tuesday, 4 February 14

(Martins and Shellard)

Nambu Model



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$$p = \frac{\mu}{L^2},$$

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(Martins and Shellard)



VOS Model



For power law expansion $a(t) \propto t^{\beta}$, find attractor solution with scaling

$$L = \epsilon t, \qquad \epsilon = \sqrt{\frac{\tilde{k}(\tilde{k} + \tilde{c})}{4\beta(1 - \beta)'}}, \qquad v = \sqrt{\frac{\tilde{k}(1 - \beta)}{\beta(\tilde{k} + \tilde{c})'}},$$

More complicated VOS models can be constructed for superstrings, networks with junctions etc.

Will use comoving correlation length

$$l = \frac{L}{a} = \xi \tau,$$

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VOS Parameters



- Simulations find network parameters vary between matter and radiation eras
- Density in radiation era is greater than in matter era
- Simulations of Martins and Shellard find

 $\xi_{\rm rad} = 0.13 \longrightarrow \xi_{\rm mat} = 0.21 \qquad v_{\rm rad} = 0.65 \longrightarrow v_{\rm mat} = 0.60$

Ringeval et al find

 $\xi_{\rm rad} = 0.16 \longrightarrow \xi_{\rm mat} = 0.19$

- Strings also have small scale structure, or 'wiggliness' (Carter 2000)
- Effective coarse grained energy momentum tensor gives rescaled mass per unit length $U = \alpha \mu$,

Estimated to be

 $\alpha_{\rm rad} = 1.5 \longrightarrow \alpha_{\rm mat} = 1.9$ Health warning: CMBACT v3



Abelian Higgs



- Lagrangian $\mathcal{L} = (D_{\mu}\Phi)^2 \frac{1}{4}F_{\mu\nu}F^{\mu\nu} V(\Phi)$
- Solve equations of motion on 3D lattice (Bevis et al)
- Extract unequal time correlator (UETC) of energy momentum tensor $\langle \Theta(k, \tau_1) \Theta(k, \tau_2) \rangle$
- Radiation into propagating modes included
- However, simulation box sizes limited to ~ 300x string core width
- Requires slowing down rate of growth of core (pre 2014, s<1)</p>
- Found only very small loop production most of energy loss to radiation
- See no observed small scale structure
- Similar network parameters in matter and radiation eras

 $\xi = 0.3 \qquad v = 0.5 \qquad \text{(Hindmarsh et al 2009)}$ More work required to resolve differences with Nambu model Cesa

CMB Anisotropies

- **Key** ingredient is the UETC
- Use stiff source approximation

String energy momentum tensor

 $G_{\mu\nu} = 8\pi G T_{\mu\nu} \qquad \qquad \delta G_{\mu\nu} = 8\pi G \left(\delta T_{\mu\nu} + \theta_{\mu\nu} \right)$

- Can estimate UETC directly from simulations, and use as sources in CMB codes
- For thin strings there is a useful intermediate framework called Unconnected Segment Model (USM) (Vincent et al, Albrecht et al, Pogosian and Vachaspati)
- Model strings are ensemble of uncorrelated straight segments, each moving with random velocity
- Inputs are correlation length, velocity and wiggliness, as measured from simulations





Mimic Model

- Mimic model computed with USM, choosing parameters to more closely resemble Abelian-Higgs
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- Turn off evolution of network parameters between matter and radiation eras

• Choose
$$\xi = 0.35$$
 $v = 0.4$ $\alpha = 1.05$

- Encouragingly find closer agreement between spectra
- Remember Nambu suffers less of an issue with dynamical range



- Abelian-Higgs has advantage of including radiation
- Even with simple physics USM does a good job of fitting both!



Health warning: CMBACT v3

planck





Power Spectrum



- Note: Amplitude set by mass per unit length $C_{\ell}^{\text{string}} \propto (G\mu)^2$
- Two primary contributions to spectrum density perturbations at last scattering, Kaiser-Stebbins effect along line of sight
- Peak position set by correlation length at last scattering
- Why the difference between Nambu and Abelian-Higgs?
- Normalisation determined by correlation length
- Nambu model has more small scale power smaller string correlation length in radiation era, by factor of ~1.6
- Different split into scalar, vector and tensors small scale structure of strings



AH



USM



- Let's now revisit the USM. In existing codes (CMBACT) the UETC is not computed
- Rather an ensemble of source histories are created, then averaged to find power spectra
- The EM tensor of a straight string segment is

$$\Theta_{00} = \frac{\mu\alpha}{\sqrt{1-v^2}} \frac{\sin(k\hat{X}_3\xi\tau/2)}{k\hat{X}_3/2} \cos(\chi + k\hat{X}_3v\tau), \qquad \Theta_{ij} = \left[v^2\hat{X}_i\hat{X}_j - \frac{1-v^2}{\alpha}\hat{X}_i\hat{X}_j\right]\Theta_{00},$$

$$\hat{X} \text{ and } \hat{X} \text{ are randomly orientated unit vectors satisfying } \hat{X} \cdot \hat{X} = 0$$
Phase χ set by location of string

• Orientated wave vector as $\mathbf{k} = k\hat{k}_3$, perform scalar, vector, tensor split

$$\Theta^{S} = (2\Theta_{33} - \Theta_{11} - \Theta_{22})/2, \qquad \Theta^{V} = \Theta_{13}, \qquad \Theta^{T} = \Theta_{12}.$$



USM



- ► During scaling number density of strings scales as $n(\tau) \propto \tau^{-3}$, requires tracking very large number of segments
- USM consolidates all segments that decay at some time into a single segment
- Number of segments that decay between τ_i and τ_{i-1} is

$$N_d(\tau_i) = V[n(\tau_{i-1}) - n(\tau_i)],$$

EM tensor of network (K consolidated segments) is

$$\Theta_{\mu\nu} = \sum_{i=1}^{\kappa} [N_d(\tau_i)]^{1/2} \Theta^i_{\mu\nu} T^{\text{off}}(\tau, \tau_i, L_f),$$

Note consolidated segment has weight $\sqrt{N_d}$ T^{off} is segment decay function



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From these it is possible to work out the USM UETC analytically

(Avogoustidis et al 2012)

$$\langle \Theta(k, \tau_1) \Theta(k, \tau_2) \rangle = \frac{2f(\tau_1, \tau_2, \xi, L_f)}{16\pi^3} \int_0^{2\pi} d\phi \int_0^{\pi} \sin\theta d\theta \int_0^{2\pi} d\psi \int_0^{2\pi} d\chi \Theta(k, \tau_1) \Theta(k, \tau_2),$$

- The scaling function can be computed analytically from the segment decay $f(\tau_1, \tau_2, \xi, L_f \to 1) = \frac{1}{[\xi \operatorname{Max}(\tau_1, \tau_2)]^3}.$
- ► UETC involves several integrals which are doable, e.g.

$$\frac{1}{2} \int_0^{\pi} d\theta \sin^3\theta \cos(x\cos\theta) J_0(\rho\sin\theta) = \left[1 + \frac{\partial^2}{\partial x^2}\right] \left(\frac{\sin\sqrt{\rho^2 + x^2}}{\sqrt{\rho^2 + x^2}}\right),$$

And only two which weren't (2), e.g.

$$\frac{1}{2} \int_0^{\pi} d\theta \sin\theta \cos(x \cos\theta) J_0(\rho \sin\theta) \sec^2\theta = \sum_{c=0}^{\infty} \frac{1}{c!} \frac{\rho}{(2c-1)} \left(-\frac{x^2}{2\rho}\right)^c j_{c-1}(\rho),$$



Comparison





Eigenmodes



UETC can be decomposed into eigenmodes - each of these are coherent, can be used as source functions in CMB code

(Pen et al 1997)

$$(k^{2}\tau_{1}\tau_{2})^{\gamma}(\tau_{1}\tau_{2})^{1/2}\langle\Theta(k,\tau_{1})\Theta(k,\tau_{2})\rangle = \sum_{i=1}^{n}\lambda_{i}u_{i}(k\tau_{1})\otimes u_{i}(k\tau_{2}),$$

Diagonalization introduces a change of basis, but since modes are orthogonal and perturbations are linear, CMB sources are

$$\Theta(k\tau) \rightarrow \frac{u(k\tau)}{(k\tau)^{\gamma} \tau^{1/2}}.$$

Power spectra found by summing over eigenmodes (ordered from highest to lowest) and truncated at some number

$$C_{\ell} = \sum_{i=1}^{n} \lambda_i C_{\ell}^i,$$



CMB Spectra



- CMB spectra are $C_{\ell}^{i(I)} = \frac{2}{\pi} \int k^2 dk \Delta_{\ell}^{i(I)}(k, \tau_0) \Delta_{\ell}^{i(I)}(k, \tau_0)$,
- Incorporated into CAMB CAMBACT

(Avogoustidis et al 2012)

E.g. scalar spectra





Some Numbers



- Running CMBACT for 2000 source realisations takes ~ 20 hours
- Impossible to run any sort of MCMC with Planck (other than fitting for the overall amplitude of a given spectrum)
- With CAMBACT we need somewhere between 50-100 eigenmodes for reasonable accuracy
- Each mode requires running CAMB for scalars, vectors and tensors (i.e. 150-300 CAMB evaluations)
- CAMB takes ~ 1 second to run, hence total CAMBACT computational time is several minutes
- A big improvement, Planck MCMC implementation work in progress

Philosophy - use USM, marginalize over string parameters



Abelian Higgs



- So far focused on USM case. Idea is to use simulations (both Nambu and Abelian-Higgs) to inform network parameters
- State of the art Abelian-Higgs simulations have also been used for the Planck analysis (Bevis et al)
- Fields evolved on 1024³ grid, starting from random initial conditions designed to mimic a phase transition
- Brief diffusive period ensures system rapidly reaches scaling
- String cores are partially fattened to enlarge dynamical range
- Various runs performed to check results insensitive to string fattening parameter
- UETC's calculated at regular intervals and used in CMBEASY code



Planck Constra

Additional parameter added to LCDM, fraction of strings at $\ell = 10$

0.00

0.01

0.02

0.03

0.04

 f_{10}

0.05



Use same dataset/priors as in main cosmology paper

Defect type	Pla f_{10}	anck+WP Gu/c ²	$Planck$ f_{10}	x+WP+highL $G\mu/c^2$
NAMBU	0.015 0.033 0.028 0.043 0.055	1.5×10^{-7} 3.6×10^{-7} 3.2×10^{-7} 11.0×10^{-7} 10.6×10^{-7}	0.010 0.034 0.024 0.041 0.054	$ \begin{array}{r} 1.3 \times 10^{-7} \\ 3.7 \times 10^{-7} \\ 3.0 \times 10^{-7} \\ 10.7 \times 10^{-7} \\ 10.5 \times 10^{-7} \end{array} $

Limits on string tension significantly improved, e.g. in Nambu model string fraction <1%</p>



AH

0.06

0.07

 $P/P_{\rm max}$

Planck + WP + highL



Planck Constraints

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Strings exhibit no significance correlations with any other cosmological parameter





Conclusions



- Planck has provided improved limits on strings
- Power spectrum gives

 $G\mu/c^2 < 1.5 \times 10^{-7}$, $f_{10} < 0.015$, (Nambu) $G\mu_{\rm AH}/c^2 < 3.2 \times 10^{-7}$, $f_{10} < 0.028$. (Abelian-Higgs)

- String evolution is crucial for understanding CMB anisotropies
- Still some uncertainty in modelling approaches main problem is dynamical range
- Improved USM model will enable MCMC search of string parameter space
- ► Is the best solution to use USM and marginalize?







