

Analytic Modelling of Cosmic String & Superstring Networks

Tasos Avgoustidis
(University of Nottingham)

Based on:

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0912.4004, 0705.3395,
astro-ph/0512582, hep-ph/0410349

with: A. Achucarro, E. Copeland, M. Leite, A. Lopez, C. Martins, A. Moss, A. Nunes,
L. Pogosian, A. Pourtsidou, P. Shellard, D. Skliros, D. Steer, J. Urestilla

In progress w/ Achucarro, Leite, Lopez, Martins, Nunes & Urestilla
Charnock, Copeland, Moss & Skliros
Pourtsidou & Sakellariadou

Motivation

- **Generic** in a wide range of models (Majumdar & Davis 2002, Burgess et al 2001, Sarangi & Tye 2002, Jeannerot et al 2003)
- A number of **potential observational effects**
- Powerful **tool for probing/constraining High-Energy Physics:**

Field Theory Strings

Energy scale of
symmetry breaking

Cosmic Superstrings

String coupling
Compactification scale(s)
Warping scale

Important to understand
their cosmological evolution

The VOS model

Kibble 1985

Martins & Shellard 1996/2000

Nambu-Goto action $S = -\mu \int \sqrt{-\gamma} d^2\zeta$, $\gamma_{\alpha\beta} = g_{\mu\nu} \partial_\alpha x^\mu \partial_\beta x^\nu$

Equation of motion $\nabla^2 x^\mu + \Gamma_{\nu\lambda}^\mu \gamma^{\alpha\beta} \partial_\alpha x^\nu \partial_\beta x^\lambda = 0$

E-M tensor $T^{\mu\nu} = \frac{1}{\sqrt{-g}} \mu \int d^2\zeta \sqrt{-\gamma} \gamma^{\alpha\beta} \partial_\alpha x^\mu \partial_\beta x^\nu \delta^{(4)}(x^\lambda - x^\lambda(\zeta^\alpha))$

From E-M tensor define energy: $E(\tau) = \int_{t=\text{const}} \sqrt{h} n_\mu n_\nu T^{\mu\nu} d^3\mathbf{x} = a(\tau) \mu \int \epsilon d\sigma$

For a network define: $v^2 = \left\langle \frac{d\mathbf{x}^2}{d\tau} \right\rangle \equiv \frac{\int \frac{d\mathbf{x}^2}{d\tau} \epsilon d\zeta}{\int \epsilon d\zeta}$ $\rho = \frac{\mu}{L^2}$

Differentiate and use eom, find:

$$\dot{\rho} = -2 \frac{\dot{a}}{a} (1 - 2v^2) \rho$$
$$\dot{v} = (1 - v^2) \left(\frac{k}{R} - 2 \frac{\dot{a}}{a} v \right)$$

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Differentiate and use eom, find:

$$\dot{\rho} = -2 \frac{\dot{a}}{a} (1 - 2v^2) \rho - \frac{\tilde{c} v \rho}{L}$$

$$\dot{v} = (1 - v^2) \left(\frac{k}{R} - 2 \frac{\dot{a}}{a} v \right)$$

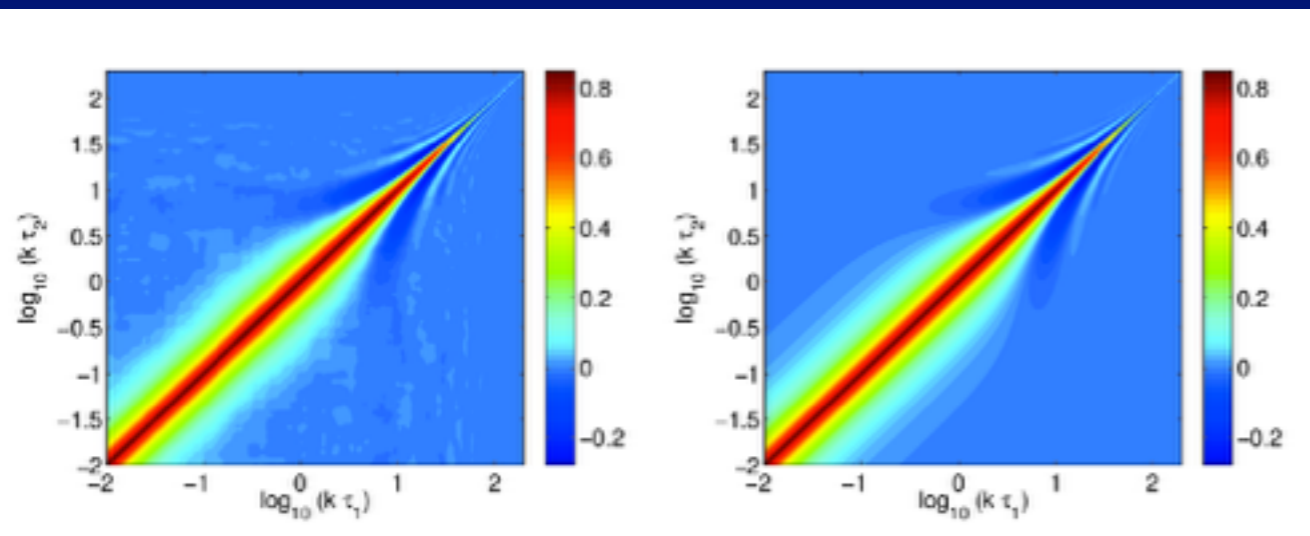
Cf simulations

Computing CMB signals from strings

- Strings are **Active, Incoherent** sources \longrightarrow need **UETC**:

$$\langle \Theta(k, \tau_1) \Theta(k, \tau_2) \rangle = \frac{2f(\tau_1, \tau_2, \xi, L_f)}{16\pi^3} \int_0^{2\pi} d\phi \int_0^\pi \sin \theta d\theta \int_0^{2\pi} d\psi \int_0^{2\pi} d\chi \Theta(k, \tau_1) \Theta(k, \tau_2)$$

- Standard approach: model network as **K unconnected segments** with lengths and velocities given by **VOS model**
- We have computed integrals analytically: (AA, Copeland, Moss & Skliros, 2012)



USM (~8 hours)

Our analytic result
(~20 seconds)

Get **CI's** in a **few mins**:
MCMC analysis including network
parameters now possible

(cf Adam's talk, Andrei's & Paul's talks)

Cosmic Superstrings

To describe cosmic superstring networks one must include the following **additional effects**:

- Strings Evolve in Extra Dimensions (ED)
- Reduced Intercommuting Probabilities (Quantum Interactions & ED)
- Multi-tension string components
- New Interactions: Junction formation
- Non-trivial Kinematic Constraints



VOS in Extra Dims: Microscopics

(AA & Shellard 2004)

- Metric

$$ds^2 = N(t)^2 dt^2 - a(t)^2 d\mathbf{x}^2 - b(t)^2 d\mathbf{l}^2$$

Nambu-Goto Action

$$S = -\mu \int \sqrt{-\gamma} d^2\zeta$$

Equations of motion

$$\dot{\epsilon} = -N^{-2}\epsilon \left\{ N\dot{N} + a\dot{a} \left[\dot{\mathbf{x}}^2 - \left(\frac{\mathbf{x}'}{\epsilon} \right)^2 \right] + b\dot{b} \left[\dot{\mathbf{l}}^2 - \left(\frac{\mathbf{l}'}{\epsilon} \right)^2 \right] \right\}$$

$$\ddot{\mathbf{x}} + \left\{ \frac{2\dot{a}}{a} - N^{-2} \left\{ N\dot{N} + a\dot{a} \left[\dot{\mathbf{x}}^2 - \left(\frac{\mathbf{x}'}{\epsilon} \right)^2 \right] + b\dot{b} \left[\dot{\mathbf{l}}^2 - \left(\frac{\mathbf{l}'}{\epsilon} \right)^2 \right] \right\} \right\} \dot{\mathbf{x}} = \left(\frac{\mathbf{x}'}{\epsilon} \right)' \epsilon^{-1}$$

$$\ddot{\mathbf{l}} + \left\{ \frac{2\dot{b}}{b} - N^{-2} \left\{ N\dot{N} + a\dot{a} \left[\dot{\mathbf{x}}^2 - \left(\frac{\mathbf{x}'}{\epsilon} \right)^2 \right] + b\dot{b} \left[\dot{\mathbf{l}}^2 - \left(\frac{\mathbf{l}'}{\epsilon} \right)^2 \right] \right\} \right\} \dot{\mathbf{l}} = \left(\frac{\mathbf{l}'}{\epsilon} \right)' \epsilon^{-1}$$

Energy-Momentum Tensor & Energy:

$$T^{\mu\nu} = \frac{1}{Na^3b^{D-3}} \mu \int d\zeta (\epsilon \dot{x}^\mu \dot{x}^\nu - \epsilon^{-1} x'^{\mu} x'^{\nu}) \delta^{(D)}(\mathbf{x} - \mathbf{x}(\zeta, t), \mathbf{l} - \mathbf{l}(\zeta, t))$$

$$E = \int_{t=\text{const}} \sqrt{h} n_\mu n_\nu T^{\mu\nu} d^3\mathbf{x} d^{D-3}\mathbf{l}$$

VOS in Extra Dims: Macroscopics

- Energy Density (equiv. correlation length L) Intercommuting Prob

$$2 \frac{dL}{dt} = [(2 + w_\ell^2) + (2 - w_\ell^2) v_x^2 + (1 - w_\ell^2) v_\ell^2] HL + cP_{\text{eff}} v_x$$

Modify 3D terms

New term due to extra dimensional velocities

String Velocities

$$v_x \frac{dv_x}{dt} = \frac{k_x v_x}{R} (1 - v^2) - (2 - w_\ell^2) H v_x^2 (1 - v^2) - H v_x^2 v_\ell^2$$

$$v_\ell \frac{dv_\ell}{dt} = \frac{k_\ell v_\ell}{R} (1 - v^2) - (1 - w_\ell^2) H v_\ell^2 (1 - v^2) + H v_\ell^2 v_x^2$$

Effective 3D string motion slows down due to extra dimensional velocities

$$v^2 \equiv v_x^2 + v_\ell^2 \leq 1/2$$

(AA & Shellard 2004)

Intercommutation Probability (Shellard 1987)

- Jackson, Jones & Polchinski 2004 $10^{-3} < P < 10^{-1}$

- Introduce $P < 1$ in 1-scale model:

$$\rho \simeq -2H\rho - \rho P/L \quad \rightarrow \quad \rho \propto P^{-2}$$

(Jones, Stoica & Tye 2003)

Effect of Small-Scale Structure? Need **simulations**.

Flat space simulations suggest: $\rho \propto P^{-1}$

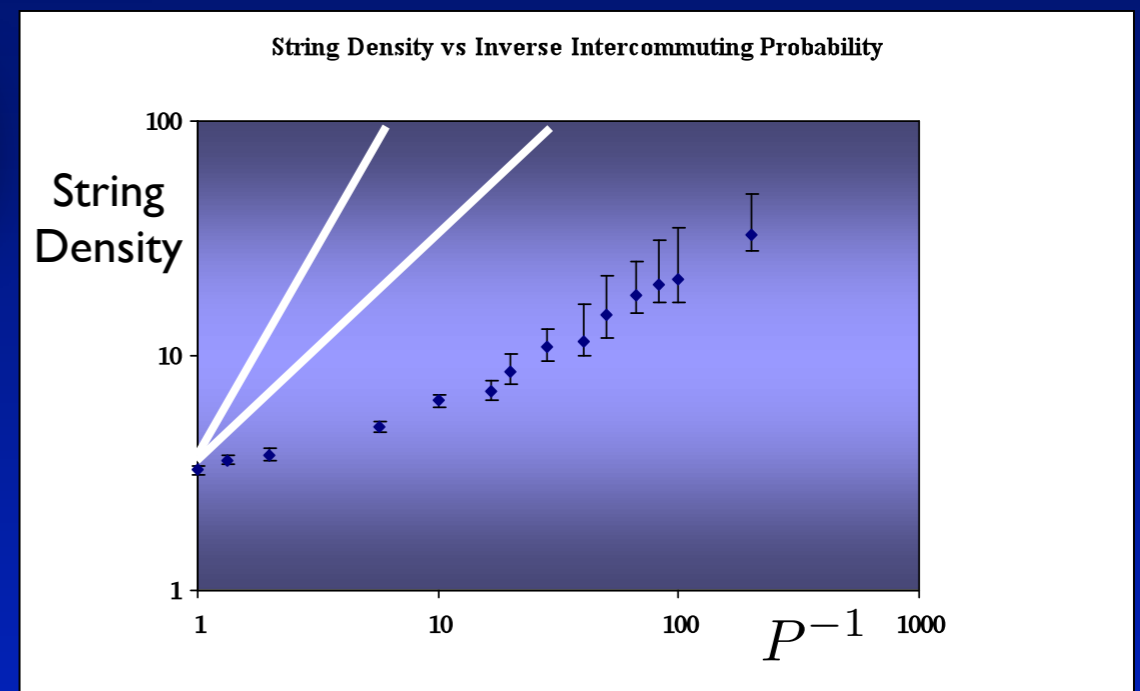
(Sakellariadou 2004)

Expanding space:
(AA & Shellard 2005)

$$\rho \propto P^{-0.6 \pm 0.1}$$

$$P_{\text{eff}} \propto P^{1/3}$$

Cf Olum & Vanchurin 2005
Vanchurin 2007



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Modelling Superstring Networks


(AA & Shellard 2007)

(cf Tye et al 2006)

For superstring networks, must include:

- N types of string (F, D, FD,...)
- junction formation
- intercommuting probs

• String Densities:



$$\dot{\rho}_i = -2\frac{\dot{a}}{a}(1 - 2v_i^2)\rho_i - \frac{\tilde{c}_i v_i \rho_i}{L_i} - \sum_{a,k} \frac{\tilde{d}_{ia}^k \bar{v}_{ia} \mu_i \ell_{ia}^k(t)}{L_a^2 L_i^2} + \sum_{b,a \leq b} \frac{\tilde{d}_{ab}^i \bar{v}_{ab} \mu_i \ell_{ab}^i(t)}{L_a^2 L_b^2}$$

Model dep params: depend on P_{ii}

Model dep params: depend on P_{ij}

String Velocities:

$$\dot{v}_i = (1 - v_i^2) \left[\frac{k_i}{R_i} - 2\frac{\dot{a}}{a}v_i + \sum_{b,a \leq b} \tilde{b}_{ab}^i \frac{\bar{v}_{ab} (\mu_a + \mu_b - \mu_i) \ell_{ab}^i(t) L_i^2}{v_i \mu_i L_a^2 L_b^2} \right]$$

Energy conservation at junctions

Length of zippers $\ell_{ij}^k(t)$ are model dependent $f(L_i, v_i, \mu_i)$

Modelling Superstring Networks

Interaction probs can be computed:

(Jackson, Jones & Polchinski 2004
Hanany & Hashimoto 2005)

Interaction (ij)	\mathcal{P}_{ij}
F-F	$g_s^2 \frac{(1 - \cos \theta \sqrt{1 - v^2})^2}{8 \sin \theta v \sqrt{1 - v^2}}$
F-D	$g_s \frac{v^2 + (\cos \theta \sqrt{1 - v^2})^2}{8 \sin \theta v \sqrt{1 - v^2}}$
F- (p, q) , $q \geq 1$	$g_s \frac{q^2 v^2 + (g_s p - \cos \theta \sqrt{(1 - v^2)(g_s^2 p^2 + q^2)})^2}{8 \sin \theta v \sqrt{(1 - v^2)(g_s^2 p^2 + q^2)}}$
D-D	$\min \left\{ \frac{\sqrt{g_s}}{2\pi^{3/4}\theta^{3/4}} e^{2\sqrt{2/3}(\theta/v)} \exp \left[-\frac{4\sqrt{\pi}\theta^{3/2}}{g_s} e^{-4\sqrt{2/3}(\theta/v)} \right], 1 \right\}$
(p, q) - (p', q') , $q, q' \geq 1$	$1 - (1 - \mathcal{P}_{DD})^{qq'}$

Non-perturbative (approximations)

- Evolution Eqns:

$$\dot{\rho}_i = -2\frac{\dot{a}}{a}(1 - 2v_i^2)\rho_i - \frac{\tilde{c}_i v_i \rho_i}{L_i} - \sum_{a,k} \frac{\tilde{d}_{ia}^k \bar{v}_{ia} \mu_i \ell_{ia}^k(t)}{L_a^2 L_i^2} + \sum_{b,a \leq b} \frac{\tilde{d}_{ab}^i \bar{v}_{ab} \mu_i \ell_{ab}^i(t)}{L_a^2 L_b^2}$$

$\tilde{c}_i \propto \mathcal{P}_{ii}^{1/3}$ Numerical Evidence

$\tilde{d}_{ij}^k \propto \mathcal{P}_{ij}^{1/3}$ Need Sims

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Kinematic Constraints

Copeland, Kibble & Steer 2006-7
 Copeland, Firouzjahi, Kibble & Steer 2008

Nambu-Goto eqns for 3 segments with a **junction**:

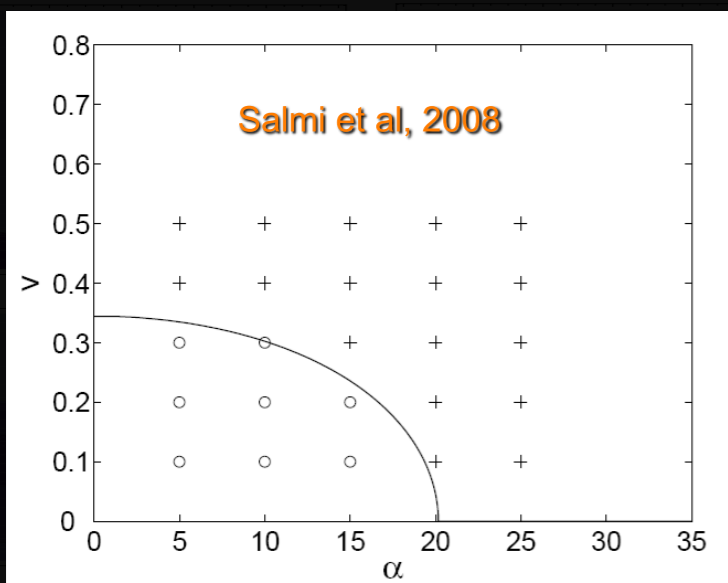
$$S = - \sum_j \mu_j \int dt \int d\sigma \theta(s_j(t) - \sigma) \sqrt{\mathbf{x}'_j{}^2 (1 - \dot{\mathbf{x}}_j^2)} + \sum_j \int dt \mathbf{f}_j(t) \cdot [\mathbf{x}_j(s_j(t), t) - \mathbf{X}(t)]$$

imply a **kinematic constraint** for junction formation:

$$f_{\vec{\mu}}(v, \alpha) < 0$$

Express as window function and integrate in (v, α) -space:

$$S_{ij}^k = \frac{2}{\pi} \int_0^1 \int_0^{\pi/2} \Theta(-f_{\vec{\mu}}(v, \alpha)) \exp[(v - \bar{v}_{ij})^2 / \sigma_v^2] d\alpha dv < 1$$



Leads to a **suppression** in the coefficients:

$$\tilde{d}_{ij}^k \rightarrow S_{ij}^k \tilde{d}_{ij}^k$$

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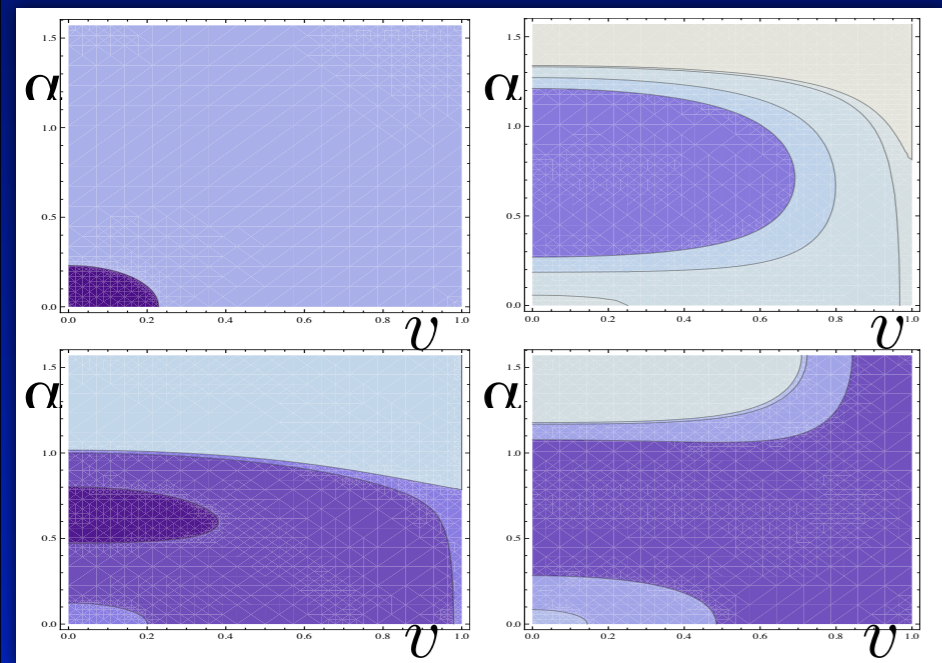
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(AA & Copeland, 2010)

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Recap

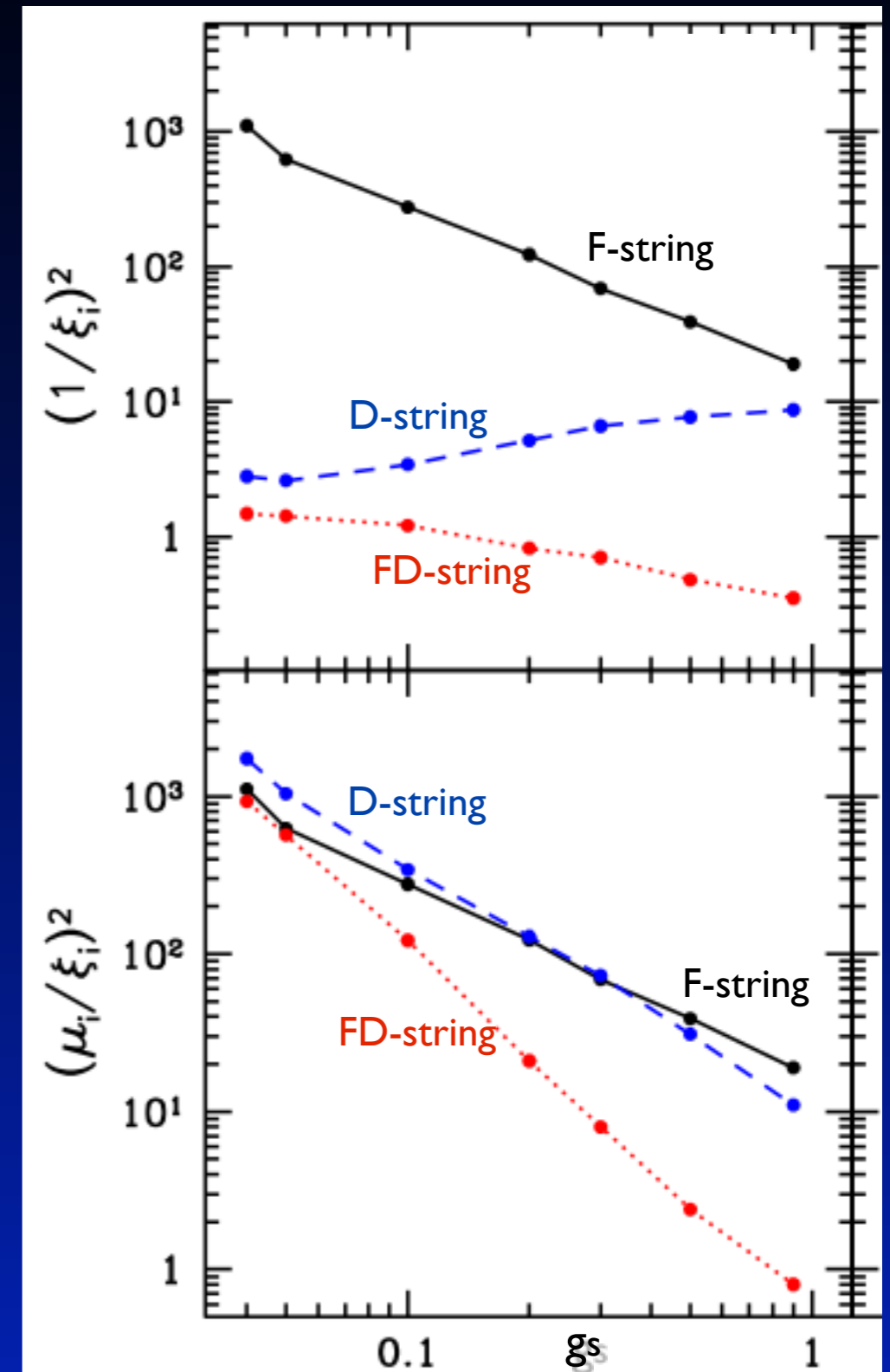
Describe interacting networks of cosmic F, D, (p,q) strings.

Started **HD** \longrightarrow Effective **3D** model with extra ingredients:

- ✓ Reduced Intercommuting Probabilities
(ED & Quantum Interactions)
- ✓ Multi-tension string components
- ✓ New Interactions: Junction formation
- ✓ Non-trivial Kinematic Constraints

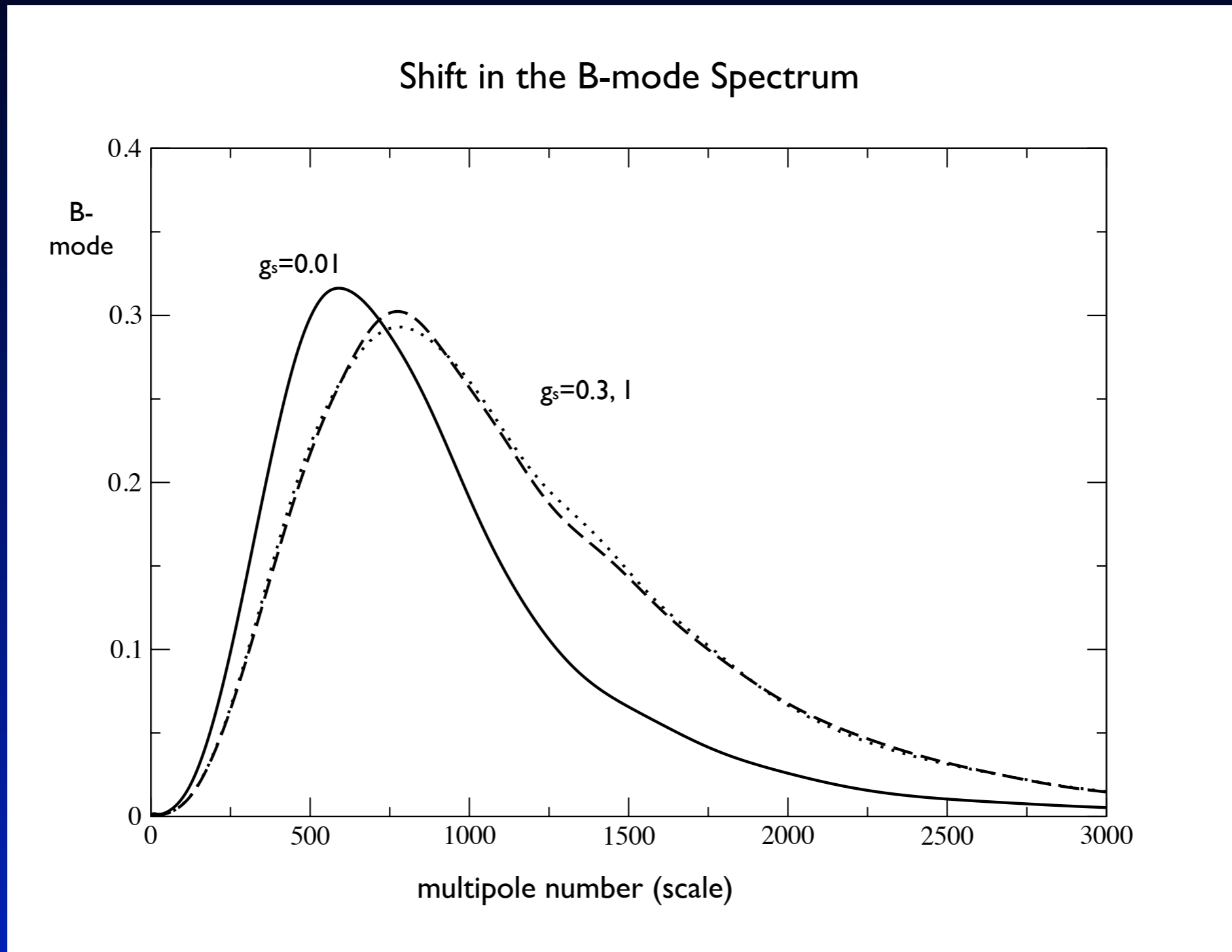
General Network Behaviour

- **Scaling for all string types**
(though we keep the first 7 lightest strings)
- **Only 3 lightest components**
(F, D, FD strings)
- **Hierarchy in number densities**
 $N_F > N_D > N_{FD}$
- **Hierarchy in tensions**
 $\mu_{FD} > \mu_D > \mu_F$
- **Number density vs “CMB” density**
Competition depending on g_s



Potentially Observable Shift in B-mode

Pourtsidou et al, 2011 (PRL)



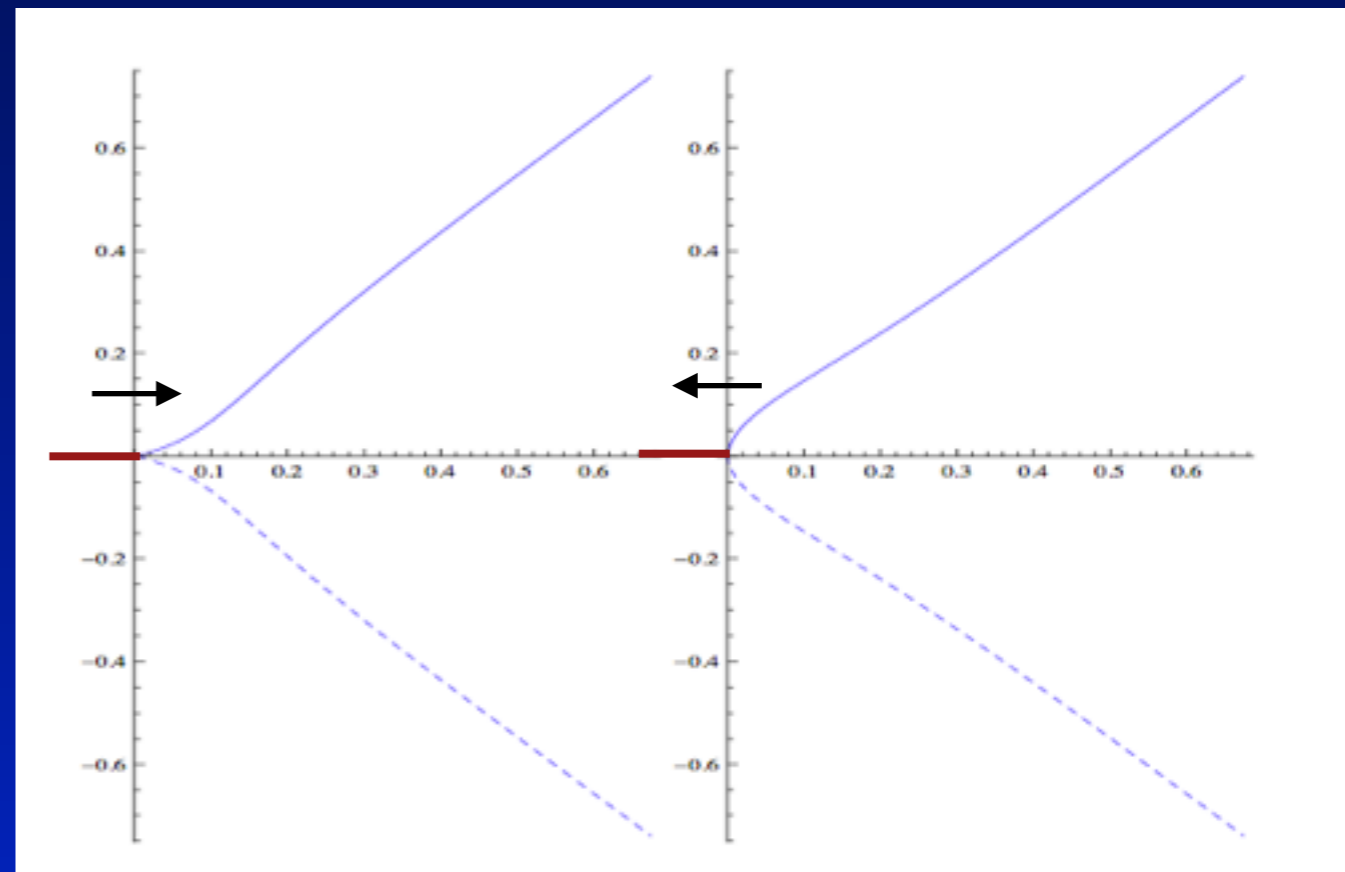
Some Problems and Open Issues

- **No simulations** to compare/calibrate with!
- **Non-perturbative amplitudes** poorly understood
- Spontaneous **unzipping???**
(In progress with Pourtsidou & Sakellariadou, see Mairi's talk)

Unzipping Mechanisms (cf Mairi's talk)

Examined a number of **potential unzipping mechanisms** in the context of Nambu-Goto strings with junctions:

- Junction Stability / Massive Monopole Dynamics **No**
- Expanding Background **No**
- String & Monopole Forces **Yes**



Semilocal Strings

Non-topological string solutions arising from the action:

$$S = \int d^4x \left[[(\partial_\mu - iA_\mu)\Phi]^2 - \frac{1}{4}F^2 - \frac{\beta}{2}(\Phi^+\Phi - 1)^2 \right]$$

(Achúcarro & Vachaspati 1999)

Strings end on monopoles: model as **hybrid string-monopole networks**:
(Martins & Achúcarro 2008)

$$3\frac{dL}{dt} = 3HL + v^2\frac{L}{\ell_d} + c_*v$$

$$\frac{dl_s}{dt} = Hl_s - v_s^2\frac{l_s}{\ell_d} + ???$$

But now have high-resolution simulations to calibrate!

(cf Ana's talk)

Semilocal Strings: Analytic Models

Model A

$$\frac{dl_s}{dt} = Hl_s - v_s^2 \frac{l_s}{l_d} + \sigma \left(1 - \frac{L}{l_s} \right) v_m^2$$

σ to be determined by simulations

Model B

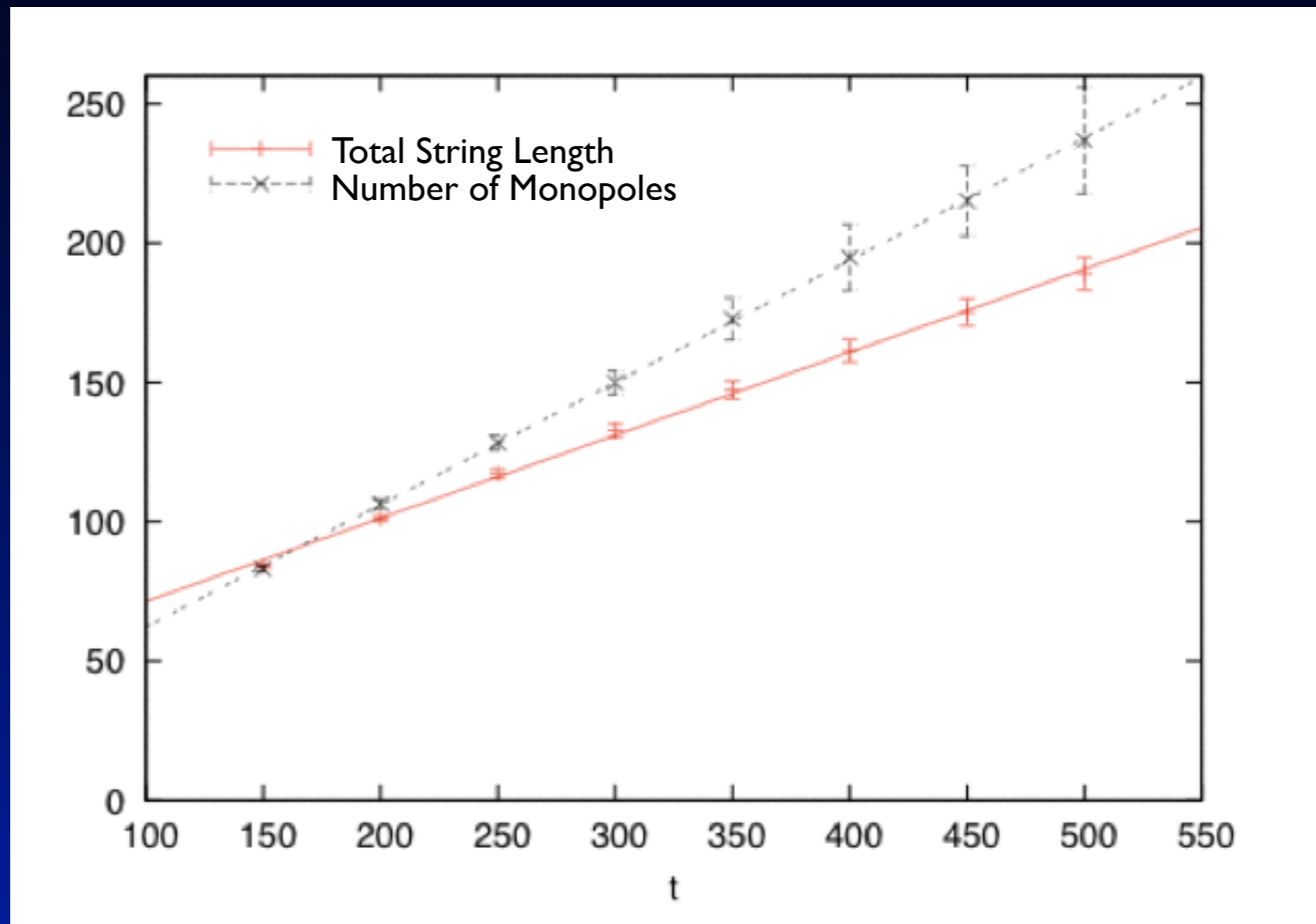
$$\frac{dl_s}{dt} = Hl_s - v_s^2 \frac{l_s}{l_d} + \left(d \frac{v_s l_s}{L} - k_1 \right)$$

d, k_1 to be determined by simulations

Semilocal Scaling

(Achucarro et al
Arxiv: 1312.2123)

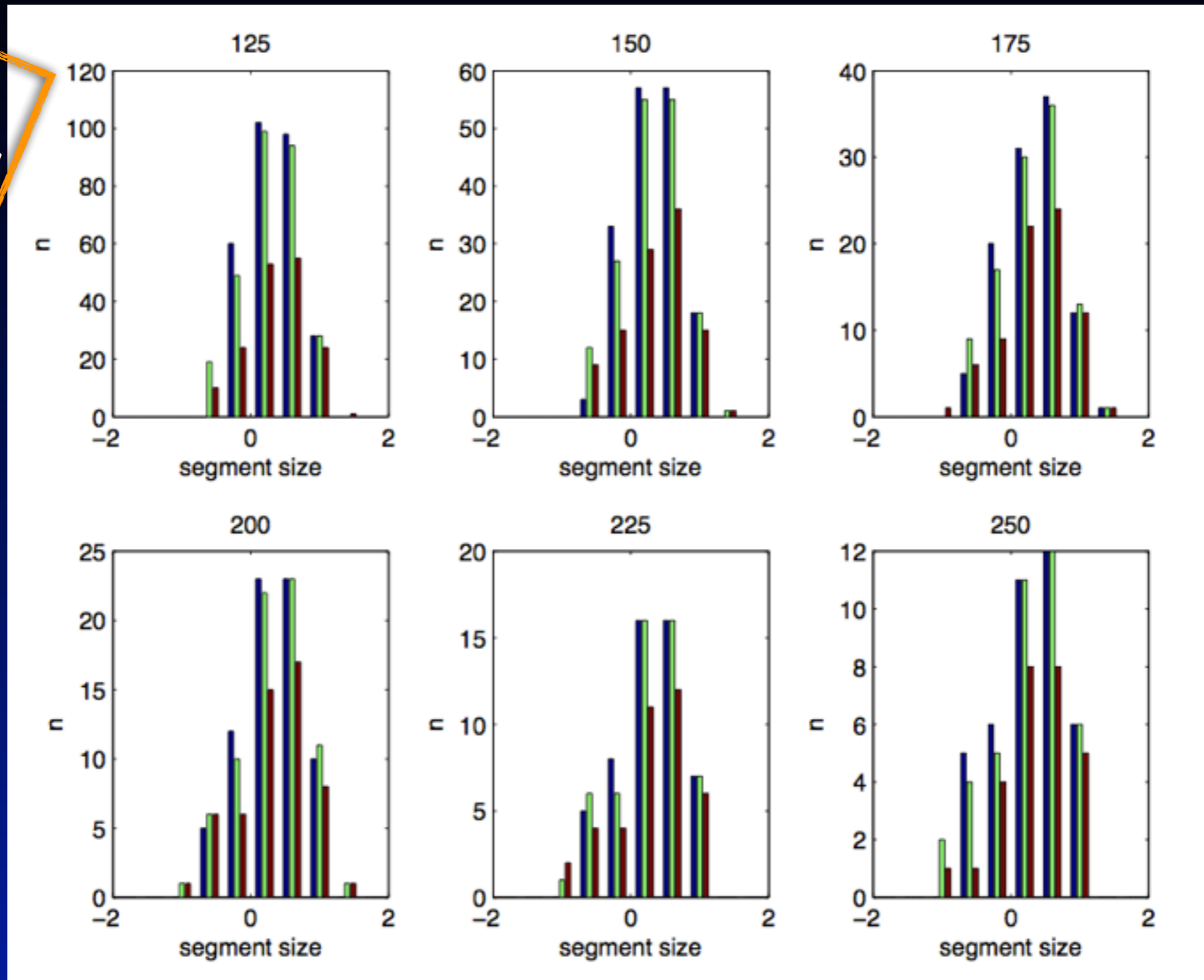
Simulations: strong evidence for **scaling both L and l_s**



Sort segments in **length bins** and compare length **distributions** of Model A, Model B, and Simulations

Semilocal Length Distributions (In progress)

Preliminary



Aim: determine dependence of model parameters on coupling β and cosmology

Conclusions & Outlook

- Field Theory Strings: analytic models demonstrated **accurate**
Currently used for **efficient CMB predictions**
- Cosmic Superstrings: Models developed but need:
 - **simulations** to test/calibrate
 - better understanding of **non-pert interactions**
- Semilocals: promising **recent progress** in both analytic & numerical modelling
 - Confronting models to sims
 - Quantifying velocities