

# Multi-Throat String Compactification and Spectrum of Cosmic String Tension

Henry Tye

HKUST and Cornell University

ASU-Tufts Joint Workshop on Cosmic Strings

Feb 3, 2014

# String theory has 10 dimensions

- 6 of the spatial dimensions are compactified;
- Flux compactification into a Calabi-Yau like manifold yields a stabilized 6-dim. internal space;
- There are D3-branes that span our uncompactified 3-dimensional space; they appear as points in the internal space;
- We may live in a stack of these D3-branes.

# A typical flux compactification in Type IIB

$$\chi(M) = 2(h^{1,1} - h^{2,1})$$

<i>Manifold</i>	$h^{1,1}$	$h^{2,1}$	$\chi$
$\mathcal{P}_{[1,1,1,6,9]}^4$	2	272	-540
$\mathcal{F}_{11}$	3	111	-216
$\mathcal{F}_{18}$	5	89	-168
$\mathcal{CP}_{[1,1,1,1,1]}^4$	1	$\mathcal{O}(100)$	$\mathcal{O}(-200)$

A manifold has  $h^{1,1}$  number of Kähler moduli and  $h^{2,1}$  number of complex structure moduli.

A typical flux compactification has hundreds of warped throats, each described by a complex structure modulus.

Example :

**P(11169)**

$$p = a_1 x_1^{18} + a_2 x_2^{18} + a_3 x_3^{18} + a_4 x_4^3 + a_5 x_5^2 \\ + \phi_1 x_1^6 x_2^6 x_3^6 + \phi_2 x_1 x_2 x_3 x_4 x_5 + \dots$$

$$x_1 \rightarrow \lambda x_1$$

$$x_2 \rightarrow \lambda x_2$$

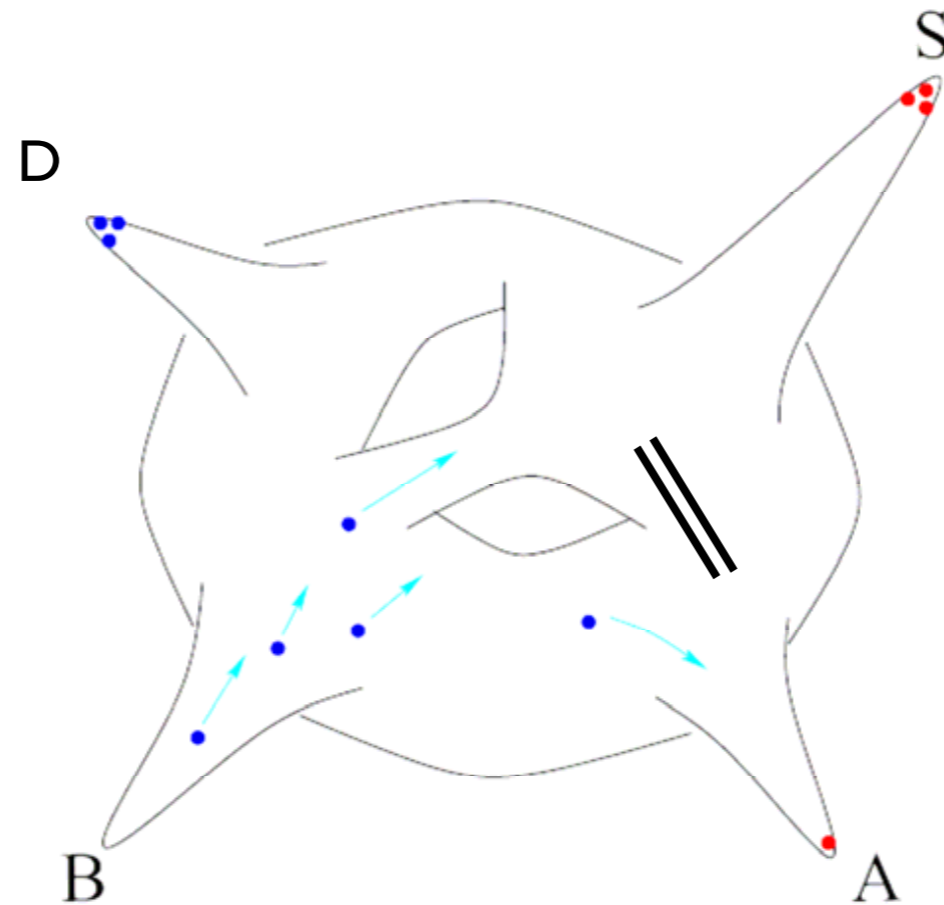
$$x_3 \rightarrow \lambda x_3$$

$$x_4 \rightarrow \lambda^6 x_4$$

$$x_5 \rightarrow \lambda^9 x_5$$

$$p \rightarrow \lambda^{18} p$$

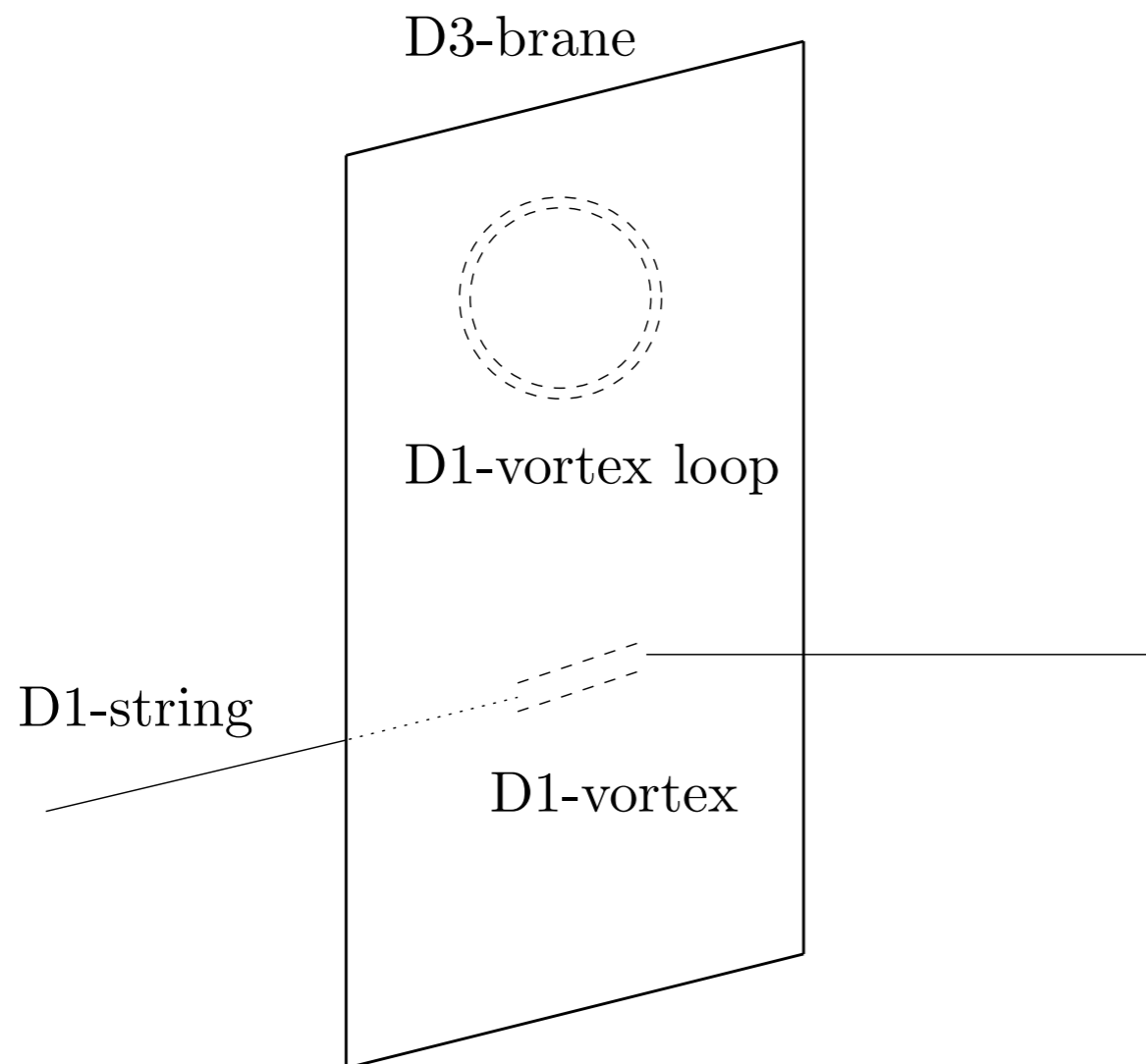
# Manifold with warped throats and D3-branes



Each warped throat has a warped factor  $h$ .

Energetically, D3-branes, D1-strings and F1 strings like to move to the bottom of the warped throats

# D1 and F1 strings



F1 strings break up into pieces inside D3-branes

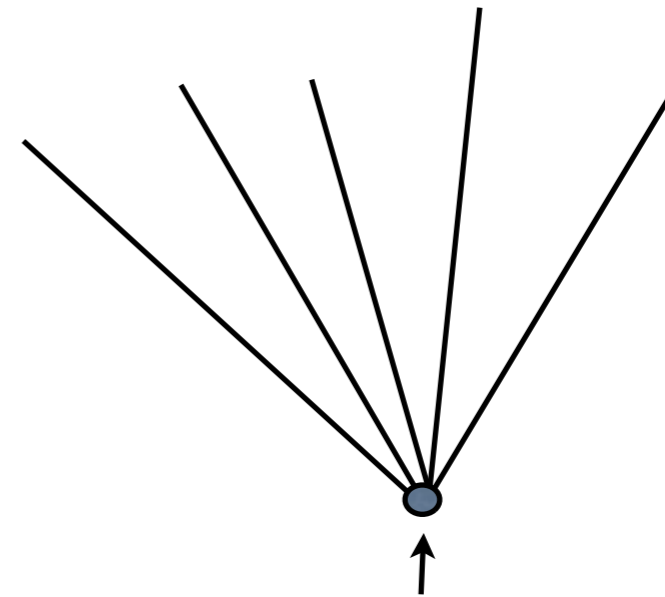
D1-strings become (metastable) vortices

Size of vortex  $r \sim \frac{\sqrt{\alpha'}}{h\sqrt{g_s}}$

# Cosmic string tension spectrum in a warped deformed throat

E.g. Klebanov-Strassler Throat :

A baryon with mass  $\sim M^{3/2} h_A / \sqrt{\alpha'}$



$$T_{p,q} \simeq \frac{h_A^2}{2\pi\alpha'} \sqrt{\frac{q^2}{g_s^2} + \left(\frac{bM}{\pi}\right)^2 \sin^2\left(\frac{\pi p}{M}\right)}, \quad b = 0.93266$$

$b=0$  in the presence  
of branes

Very large  $M$  in bulk

Leblond, Firouzjahi, HT, hep-th/0603161  
Herzog, Klebanov, hep-th/0111078

# Possibilities

- Cosmic strings may fall into the throats. Each warped throat (with no D3-branes in its bottom) has its own  $(p,q)$  spectrum.
- A throat with D3-branes at its bottom will only have D1 vortices.
- Presumably most if not all the throats have their own cosmic string networks. Each throat has its own tension scale, typically some orders of magnitude below the string scale.
- This results in high cosmic string density (i.e., sum of cosmic string networks).



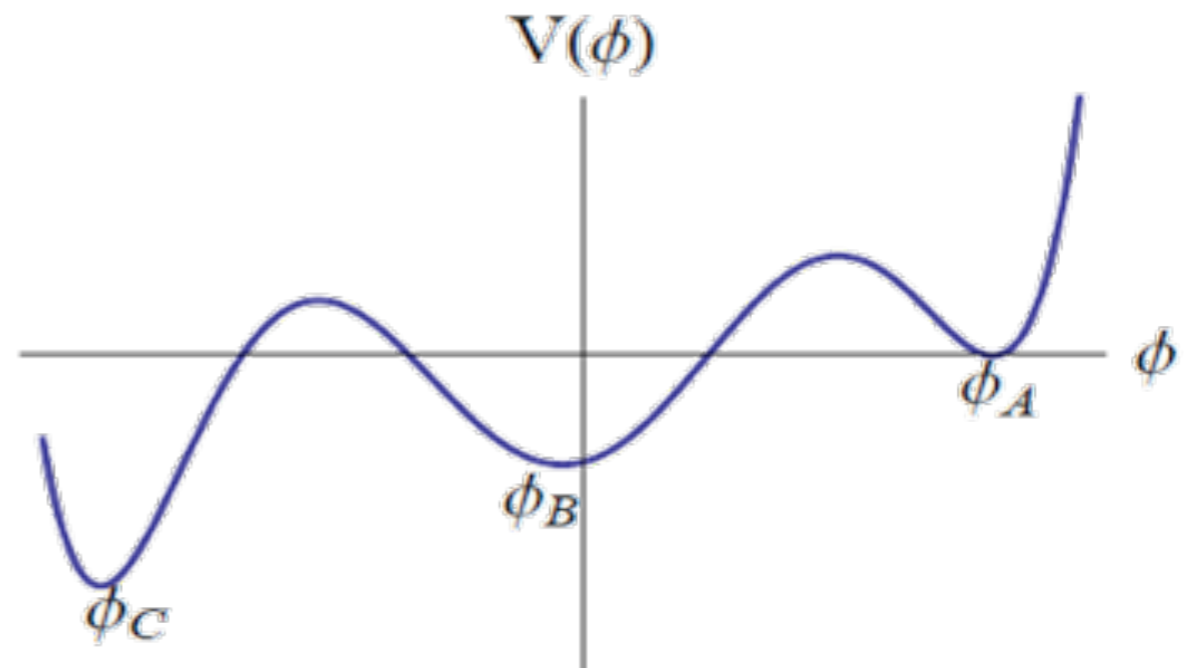
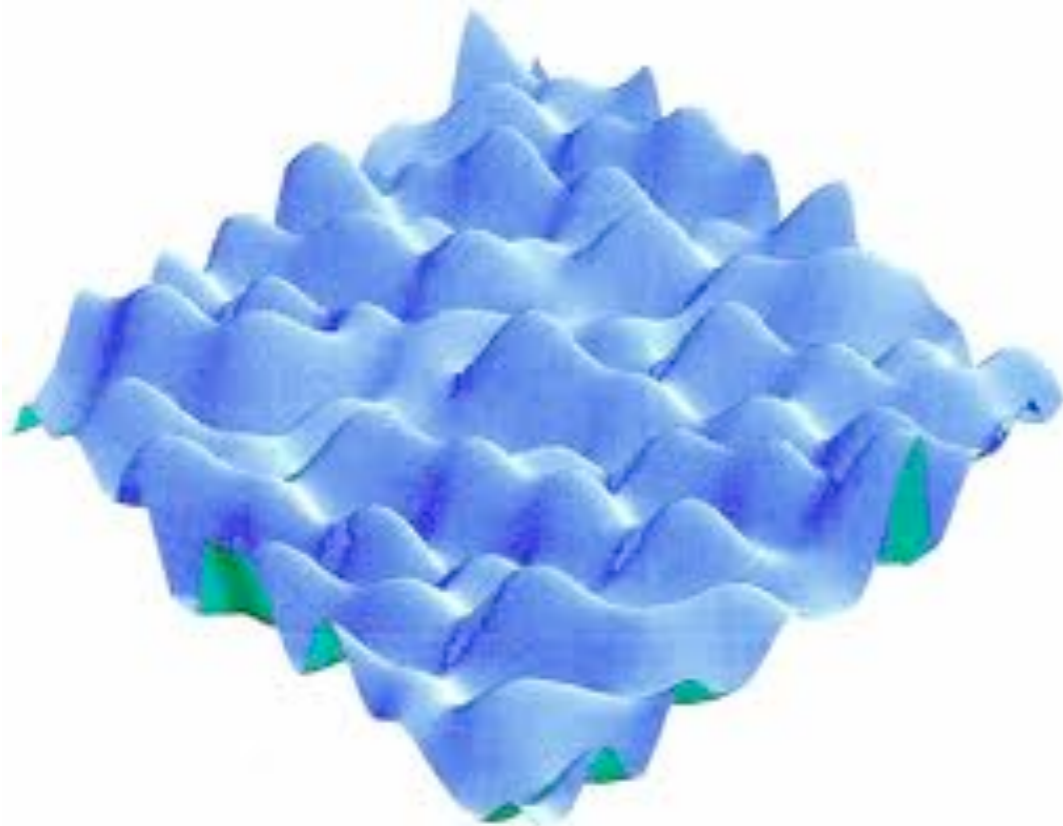
# Another possibility

- Cosmic strings may continue to oscillate in the internal space because the damping into a throat is too weak. [Avgoustidis, 0712.3224](#)  
[Avgoustidis, Chadburn, Gregory, 1204.0973](#)
- Strings moving in 9 spatial dimensions do not see each other, so intercommutation probability is small and the cosmic string density is very high.

# 3 properties to consider

- If a throat has no branes, the beads and junctions can slow down the strings.
- If a throat has D-3 branes, or D7-branes are around, the thickness of the D1-vortex
- Many throats

$$r \sim \frac{\sqrt{\alpha'}}{h\sqrt{g_s}}$$



# Summary

The cosmic string density can be much higher than expected,  
either  
as a sum of cosmic string networks,  
or  
as strings moving in internal space.

In both cases, one should detect different string tensions.